

§46. Effect of Pellet Injection Location on Ablation and Subsequent Density Redistribution

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Significant improvement in fueling efficiency has been demonstrated with the magnetic high-field-side (HFS) pellet injection[1][2] in recent tokamak study. The mechanism for the improvement is not yet elucidated but could be caused by an $E \times B$ drift that arises from a vertical polarization of ablation cloud due to a ∇B and curvature drift[3]. Noticing the difference between tokamak and LHD on the magnetic structure, pellet injection experiments from the different locations have been carried out on LHD for the purpose of investigating effect of the magnetic structure on the density redistribution.

Because the external coils, namely, a pair of helical coils establish magnetic confinement field, there are the following remarkable magnetic structures on LHD, (1)Fully negative shear and large rotational transform ($q < 1$) at the plasma boundary, (2)saddle point of the magnetic field exist in the plasma center, (3)two kind of the HFS can be defined, i.e. the toroidal averaged-HFS and the local-HFS.

The injection locations are shown in fig. 1. The OPI (outer port injection) is injection from the torus outside and equal to the LFS on tokamak. The IPI (inner port injection) is injection from the averaged-HFS. CSI (coil side injection) is injection from the local-HFS. Typical pellet speed is 1200 m/s at conventional OPI, but because the pellet speed is limited in the curved guide tubes, only slow (up to 300 m/s)[4] pellet injection is available at alternative injection. Specially modified injector with mechanical punch was employed for slow pellets acceleration.

A temporal changes of electron density and ablation light (H_{α} intensity) for the CSI and OPI are shown in fig. 2. The penetration depths, which are estimated by duration of H_{α} emission from the ablating pellet, are 0.28 m (OPI) and 0.18 m (CSI). Considering the ablation in an ergodic layer, this penetration depth is consistent with the predicted value from the NGS models. Looking at the contour plot of the magnetic field intensity on poloidal section (fig. 1), OPI and CSI are obviously different from each other in the gradient of the magnetic field. From the direction of ∇B , we have expected preferable density redistribution on CSI case. However the evident density redistribution is not observed on CSI. Though this result is quite different from the behavior in tokamaks, we tried to understand this result from the viewpoint of the $E \times B$ drift model. In consideration of the axis asymmetric field structure on LHD, $E \times B$ drift direction on the field line was calculated. For the qualitative discussion, only direction of vector was calculated, i.e. $E \times B \approx -c(\mathbf{B} \times \nabla B) \times \mathbf{B} = -\nabla B$. Fig. 3 shows a plot of the magnetic field strength versus the toroidal turn on the flux surface of $\rho=1.0, 0.7, 0.5, 0.2$. Due to helical coils, there are helical ripples and large rotational transform at the plasma boundary region where the slow pellet ablates. Therefore the

magnetic field strength alternate drastically with short period, while these effect become moderate in the core region of the plasma. The ∇B direction and poloidal angle at the point along the field line that begins at the ablation region of the slow pellet ($\rho=0.9$), is shown in fig. 4. ∇B makes 9 revolutions during one toroidal turn, i.e. direction of the ∇B polarization alternates with each 1 m moving of the pellet cloud along the field line. Therefore electric field, which drives the density redistribution, cannot be formed sufficiently in the pellet cloud. In addition, large rotational transform rapidly even out the influence of injection location on the density redistribution. Thus the density redistribution behavior on LHD can be understood qualitatively from the $E \times B$ drift model.

Reference

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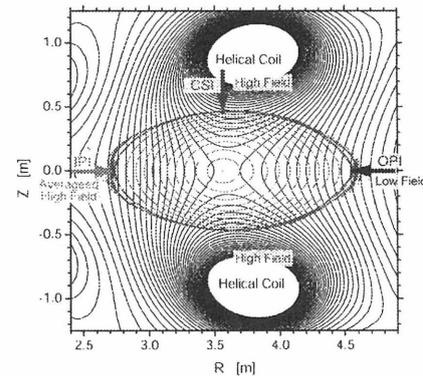


Fig. 1 Magnetic field structure and schematic view of pellet injection location in LHD.

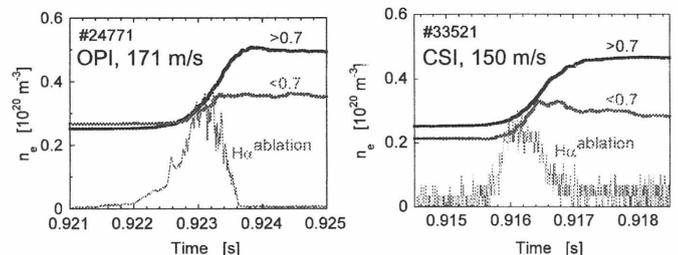


Fig. 2 Temporal changes of the electron density and H_{α} intensity at the low speed OPI and CSI.

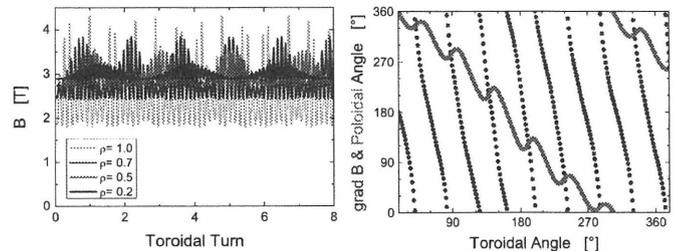


Fig. 3 Plot of the magnetic field strength versus toroidal turn on the flux surface of $\rho=1.0, 0.7, 0.5, 0.2$.

Fig. 4 Plot of the ∇B direction (same definition as poloidal angle) and poloidal angle versus toroidal angle on the flux surface of $\rho=0.9$.