

§25. Two-dimensional Density Fluctuation Measurement Using Beam Emission Spectroscopy in Heliotron J

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In order to understand the anomalous transport induced by fluctuations (MHD/turbulence) in high temperature plasmas, it is important to know spatial structure of the fluctuations experimentally. Beam emission spectroscopy (BES) is a method to measure the density fluctuation at a local position using neutral beam injection (NBI). We have installed the BES system into the Heliotron J device [1,2], which has 16 viewing chords in the radial direction. This system has a capability to measure the radial profile of the density fluctuation in the whole ($0 < r/a < 1$) plasma region in one-discharge [3]. In this FY, we have tried to investigate the propagation characteristic of the MHD fluctuation in the radial and the poloidal directions using 2×2 (radial \times poloidal) sightlines in order to discuss the feasibility of 2-dimensional BES measurement in Heliotron J.

The experiments were carried out in ECH + NBI (balanced injection, $P_{\text{NBI}}=1.2\text{MW}$) plasmas at the density of $\bar{n}_e = 0.3 \times 10^{19} \text{m}^{-3}$. The 2×2 BES sightlines observed at $r/a=0.20$ and 0.28 and $\Delta z = 0.01$ (m). As shown in Fig. 1(a), a relatively wide-band density fluctuation was observed in the frequency range of $f=5\sim 30\text{kHz}$. This fluctuation has relatively low coherence less than 0.3 to the Mirnov coil (see Fig. 1(b)). On the other hand, fluctuation with the frequency of $100\sim 110\text{kHz}$ has coherence more than 0.4, which is due to fast-ion induced MHD activity.

In order to investigate the propagation of the density fluctuation both in the radial and the poloidal direction, we applied a frequency-wavenumber analysis. The frequency-wavenumber spectrum $S(f,k)$ between two adjacent viewing chords (a and b) is determined as,

$$S_{ab}(f,k) = \frac{1}{M} \sum_{j=1}^M I_{\Delta k} [k - k^j(f)] \times \frac{|S_a^j(f) + S_b^j(f)|}{2}, \quad (1)$$

where $S(f)$ and M are the fluctuation intensity and the number of ensembles, respectively. Figure 2(a) and (b) show the frequency-wavenumber spectra for the density fluctuation in the frequency range of $5\sim 30\text{kHz}$ in the radial and the poloidal directions, respectively. In this frequency range, the wavenumber in the radial direction k_r was in the range of -0.5 to 0.5 (cm^{-1}). On the contrary, the poloidal wavenumber k_θ was around -0.5 to -1.0 (cm^{-1}), propagating in the ion-diamagnetic-drift direction, and it was not so sensitive to the frequency. This means that the phase velocity in the poloidal direction may vary with changing the frequency.

In this analysis, the turbulent broadband fluctuation has not been identified yet because of inadequate signal-to-noise ratio. To upgrade to the 2-dimensional BES measurement, improvement in the objective optical systems is required to obtain the beam emission effectively as well as to expand the observation area. Now, we are planning to install a imaging optical fiber and a 2-dimensional detector system such as ultra fast framing ($> 1\text{MHz}$) camera with high sensitivity.

- 1) S. Kobayashi, et al., Rev. Sci. Inst. **81**, 10D726 (2010).
- 2) S. Kobayashi, et al., Rev. Sci. Inst. **83**, 10D535 (2012).
- 3) S. Kobayashi, et al., 40th EPS conf. ECA vol. 37D, (2013) P1.148.

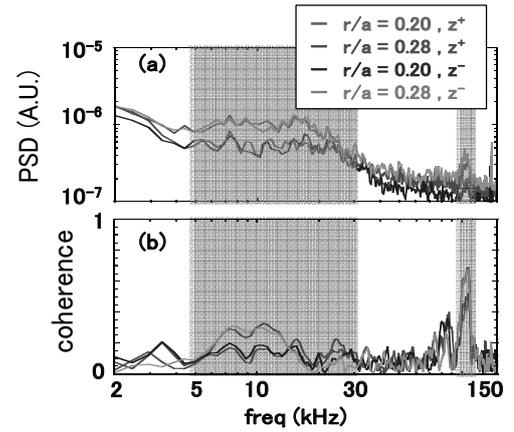


Fig. 1. (a) Power spectrum of density fluctuation measured with BES using 2×2 (radial \times poloidal) sightlines and (b) coherence to Mirnov coil.

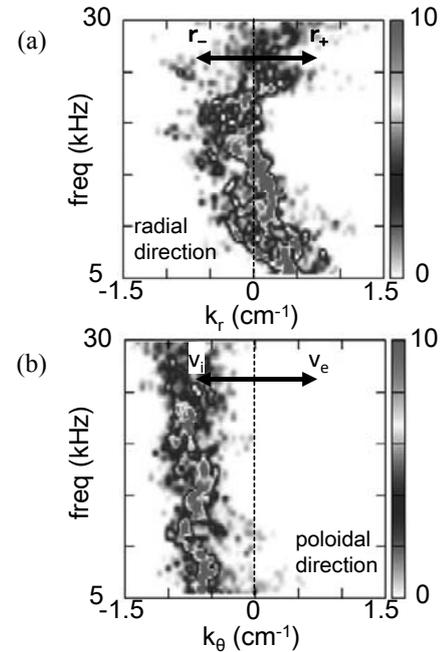


Fig. 2. Frequency-wavenumber spectra $S(f,k)$ for the density fluctuation in the frequency range of $5\sim 30\text{kHz}$ between (a) radially and (b) poloidally adjacent sightlines.