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Design Optimization of a Fast-Neutron Detector with Scintillating Fibers for Triton Burnup Experiments at Fusion Experimental Devices

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Time-resolved triton burnup studies have been carried out to estimate the behavior of alpha particles in DD fusion experimental devices. In those studies, 14 MeV neutrons emitted through DT reactions in DD plasmas should be measured selectively in the backgrounds of DD neutrons and gamma rays. For that purpose, a scintillating fiber (Sci-Fi) based fast-neutron detector has been adapted because of its advantages such as fast response, design flexibility in detection efficiency by changing the number of Sci-Fi and discrimination property against 2.4 MeV neutrons produced through DD reaction and gamma rays. However, as an optimization study of its design parameters to meet the requirements as 14 MeV neutron detector has never been done, its length had conventionally set to around 10 cm. In the present study, we tested three types of Sci-Fi detectors with three different lengths and compared with the simulated results of energy deposition, through which we tried to understand the phenomena in the detection process of fast neutrons. From the results, it has been shown that, due to the self-shielding of neutrons by Sci-Fi and the attenuation of scintillation photons during the transmission process to the photomultiplier tube, the optimal length of Sci-Fi is concluded to be about 6 cm.

I. INTRODUCTION

In deuterium operation of fusion experimental devices, to assess energetic-alpha-particle confinement performance, the behavior of 1 MeV triton produced by $d(d, p)t$ reaction has been studied. It is because the kinetic parameters of alpha particles produced by DT reaction such as Larmor radius and precession frequency are similar to those of tritons^[1]. Also, it has been shown that the initial velocity distribution of triton is isotropic as is the case of alpha particles. Time-resolved triton burnup studies have been performed by measuring secondary DT neutron created by $t(d,n)\alpha$ using scintillating-fiber (Sci-Fi) detectors in large tokamaks^[2-5] and helical system^[6-9] because of their properties such as fast response, high counting efficiency, gamma suppression ability, and directional response to incident neutron. However, in those studies, the parameters of Sci-Fi detectors

36 have been similar, namely the length of Sci-Fi were set to around 10 cm. Considering the demands for Sci-Fi detectors in
37 various fusion experimental devices as LHD and KSTAR etc. [6-9], it is required to optimize the design parameters depending
38 on the radiation fields for each experimental device. In a previous study, the authors compared the simulated and measured
39 responses of Sci-Fi detector under neutron and gamma irradiation to understand the behavior of the detector^[10]. In the present
40 study, after preparing the detectors with low transmission loss of scintillation photons, the physical phenomena occurring in
41 Sci-Fi detectors upon fast neutron measurements as energy production of recoil protons, energy deposition by protons,
42 production of scintillation photons, propagation of photons including loss have been studied through simulation and
43 experiments by using fast-neutron sources. From the results, the optimum length of Sci-Fi was revealed to be 6 cm.

44 II. Sci-Fi DETECTOR

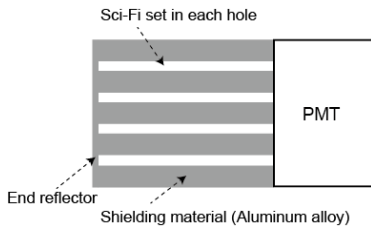
45 Sci-Fi detector was developed in the 1990s to measure 14 MeV neutron from the secondary DT reactions from the fusion
46 experimental devices such as TFTR and JT-60U [2-5]. From the measured results with 10 cm-long Sci-Fi detectors, it was
47 shown that 14 MeV neutrons could be selectively measured near the fusion devices being operated by DD discharge. In the
48 experiments at JT-60U, 14 MeV neutrons generated in the central and the peripheral regions were measured with the Sci-Fi
49 detectors. The results demonstrated the feasibility of this detector to measure time resolved triton burnup in each region, which
50 provided valuable knowledges such as toroidal Alfvén eigenmodes principally affected fast tritons in the plasma periphery.

51 As shown in Fig. 1, Sci-Fi detector is composed of scintillating optical fibers (Sci-Fi) set in each hole made through metal
52 material such as aluminum alloy. A photomultiplier tube (PMT) was set to one end and the other end is covered also with
53 aluminum alloy with reflective surface. When a fast neutron enters a Sci-Fi, a recoil proton is generated which deposits energy
54 inside the Sci-Fi core and scintillation photons are generated. The photons transmit to the ends of the Sci-Fi to reach the PMT
55 set at one end of the Sci-Fi head. If the Sci-Fi is long enough, transmitting photons are limited to those which are generated
56 for the directional angle larger angle than the critical angle of Sci-Fi. However, as Sci-Fi detector length is usually less than
57 10 cm, scintillation photons emitted to smaller angle can transmit to the PMT after being reflected at the inner surface of the
58 hole in which the Sci-Fi is set. If the hole surface is rough, the number of scintillation photons which can reach the PMT
59 depends largely on the position of generation. When a proton is recoiled to the direction nearly parallel to the Sci-Fi, it can
60 deposit large energy inside its core, however, if it is generated to the direction with larger angle to the Sci-Fi, its energy
61 deposition becomes small. Also, recoil proton is generated with larger probability to the direction close to the neutron incident
62 direction. Therefore, by setting an appropriate discrimination level to the measured pulse height by the PMT, it can realize a
63 directional property to the neutron. The authors have been trying to understand the physical phenomena occurring during the

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64 measurement process^[10] and, in this study, after improving the condition of Sci-Fi detector as polishing the hole surfaces where
65 the Sci-Fi are set, carried out the optimization of the detector configuration.

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67 FIG. 1. Structure of Sci-Fi detector.

68 II. EXPERIMENT AND SIMULATION

69 In the present study, Sci-Fi detectors were fabricated with the scintillating optical fibers: SCSF-78M (Kuraray) which were
70 connected to the PMTs: H11934-100MOD (Hamamatsu). The pulses from the PMTs were converted to digital signals with
71 the fast DAQ: DT5751 (CAEN), where the sampling was carried out at 1 GHz and the sampled values were integrated for each
72 pulse. To understand the physical phenomena, we carried out the experiments with the accelerator neutron sources,
73 OKTAVIAN^[11] at Osaka University and Fast Neutron Laboratory (FNL)^[12] at Tohoku University, and a ²⁵²Cf spontaneous
74 fission neutron source. To see the position dependencies of detection efficiency, the gamma ray sources of ¹³⁷Cs and ⁶⁰Co were
75 also used. In the experiments, Sci-Fi detector with lengths of 3 cm, 6 cm and 9 cm are used, where the number and the diameter
76 of the Sci-Fi were set to be 144 and 1 mm. As the 3 cm-long Sci-Fi detectors, we used two types, namely one which is capped
77 with aluminum reflector at the different end of the PMT and the other without it. With the latter type, we studied the effect of
78 reflection of scintillation photons by the end cap by changing the material to put at the end. Also, measurement was carried
79 out with two PMTs set at both ends of one 3 cm-long detector head to examine the possibility of efficient measurement of
80 scintillation photons.

81

82 III. RESULTS AND DISCUSSIONS

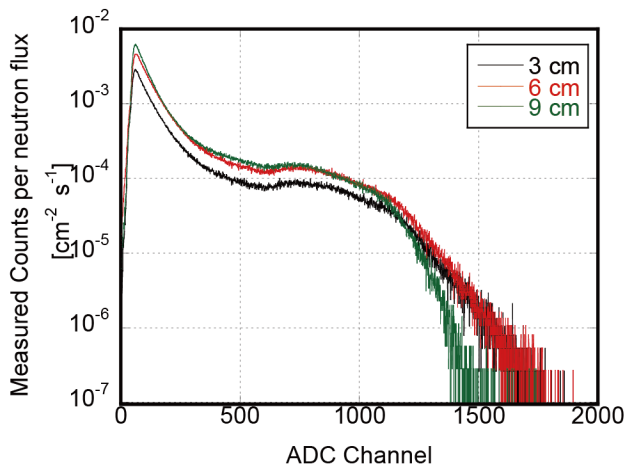
83 A. Optimization of Sci-Fi length through evaluation of measuring efficiency

84 A recoil proton is generated by a fast neutron and it deposits energy inside a Sci-Fi. Scintillation photons propagate to the
85 PMT which generates a voltage pulse. In our experimental setup, each pulse was sampled at every 1 nanosecond with the
86 resolution of 10 bits for the reference voltage of 1 V. After summing the converted values over the total pulse shape, the sum
87 was divided by 10 to set the maximum to be 1023. In Fig. 2, the measured pulse height distributions are shown for the three

88 types of detectors, namely 9 cm-, 6 cm- and 3 cm-long detectors, where 14 MeV neutrons were irradiated on them at
89 OKTAVIAN. The detectors were set at the distance of 40 cm from the Tritium target with the angle of 90 degrees from the
90 beamline. The ADC channel denotes the integrated value of the sampled values for each pulse. The data have been normalized
91 by the neutron flux at the detector position estimated with the measured results by using the ^3He proportional detector and the
92 activation foils.

93 In the highest region over 1200 ch. of the pulse height distribution shown in Fig. 2, the number of counts for the three
94 detectors were almost the same, besides the tail in the result with the 9 cm-length detector is shorter than those with the 6 cm-
95 and 3 cm-long detectors. The reason for this difference is unclear, but it may be caused by the difference of reflection condition
96 near the PMT. Although the holes in which Sci-Fi were set have been made by drilling carefully, their surface conditions may
97 be different between each detector. However, the counts over 1200 ch are less than several % of those over 500 ch. and, once
98 the detector is chosen which is to be used in fusion experimental devices, its property will be constant. We believe the influence
99 of the difference in the pulse height spectra over 1200 ch. can be ignored.

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101 FIG. 2. Measured pulse height distributions with the three types of Sci-Fi detectors, namely with the lengths of 9 cm, 6 cm,
102 and 3 cm.

103
104 In the region over 800 ch., the measured counts with the 9 cm- and 6 cm-detectors are almost comparable. This should be
105 caused by self-shielding of the neutrons by the Sci-Fi itself. Fig. 3 and Fig. 4 show the calculated distribution of recoil protons
106 near a 9 cm-long detector by the simulation code the Particle and Heavy Ion Transportation code System (PHITS) [13], where
107 14 MeV neutrons are irradiated parallel to the Sci-Fi detector. PHITS is the abbreviation of “Particle and Heavy Ion Transport

code System” which has been developed under collaboration between Japan Atomic Energy Agency, Research Organization for Information Science and Technology, and High Energy Accelerator Organization. It can deal with the transportaion of all particles over wide energy ranges, using several nuclear reaction models and nuclear data libraries. We can observe gradual decrease of the recoiled proton flux with increasing the distance from the input end which results in the comparable counts of 9 cm- and 6 cm-detectors in the highest pulse height region.

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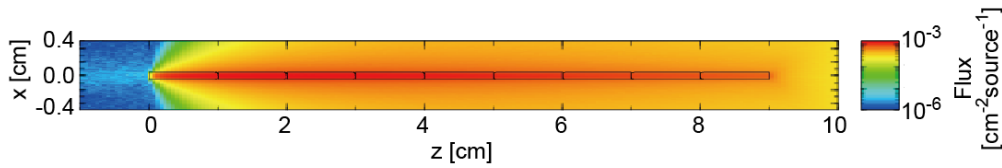


FIG. 3. Calculated recoil proton distribution when 14 MeV neutrons are input parallelly to one 9 cm-long Sci-Fi, where the intensity is shown as heat map.

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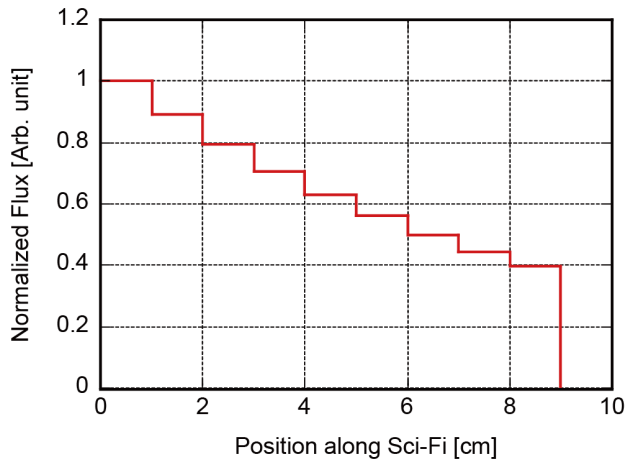


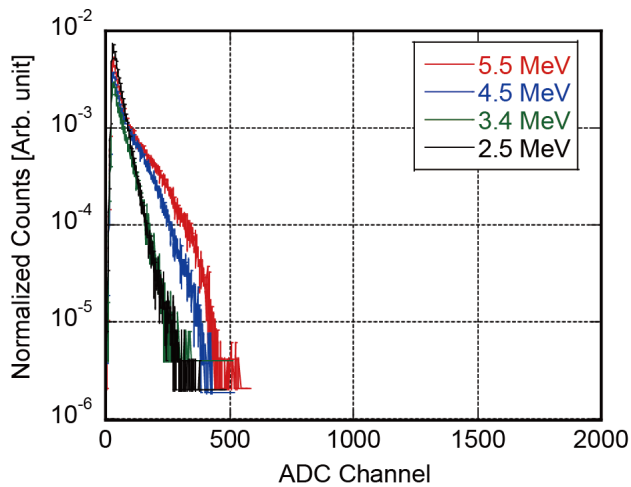
FIG. 4. Calculated dependence of normalized recoil proton flux on the distance from the input end of the scintillating fiber.

Fig. 5 shows the pulse height distributions measured at FNL with changing the neutron energy generated through DD reactions, where the counts have been normalized by the measured counts with the ³He proportional detector. We can conclude that, by setting the threshold level at around 500 ch., we can measure DT neutrons selectively in high background of DD neutrons around fusion experimental devices. Also, in Fig. 6, the results measured for ⁶⁰Co gamma rays is shown. Considering the difference of the production efficiencies of scintillation photons by proton and gamma ray, the pulse height for the gamma

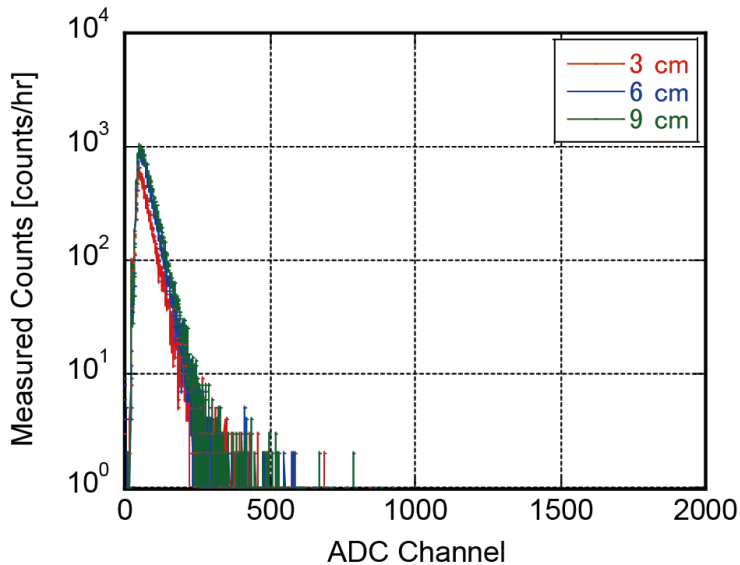
125 rays with energy of over 1 MeV corresponds to that for the protons with energy of over 2 MeV. On the other hand, from the
126 previous study on the radiation field around LHD, it is concluded that the average energy of gamma rays is 1.4 MeV^[14], which
127 means that, by setting the threshold as mentioned above, gamma ray pulses can also be discriminated. From the comparison
128 of the measured results with 9 cm- and 6 cm-long detectors shown in Fig. 3, even if we use a Sci-Fi detector longer than 6 cm,
129 the counts should be almost the same as that measured by a 6 cm-long detector. From the results, it can be concluded that the
130 optimal length of Sci-Fi detector is around 6 cm for the radiation condition around LHD.

131 The threshold for the Sci-Fi detector is determined considering the gamma ray energy at the detector position. If its energy
132 is lower than that assumed in the present estimation, the threshold can be set to the lower value, where 9 cm-long Sci-Fi
133 becomes more advantageous than the 6 cm-long Sci-Fi as more counts is observed in the lower pulse height region for 9 cm-
134 long Sci-Fi in Fig. 2. On the other hand, if the gamma ray energy is higher than the present case, as the threshold should be set
135 higher, the optimum Sci-Fi length can become shorter. Namely, 6 cm is the optimum length for the experimental devices such
136 as LHD at which the peak gamma ray energy is about 1 to 2 MeV.

137 As for the difference of 14 MeV neutron flux, the number of Sci-Fi and their diameter can be the other optimizing
138 parameters. When the 14 MeV neutron flux is low, we can enhance the efficiency by increasing the number of Sci-Fi. If we
139 choose to increase the diameter to enhance the efficiency, the pulse height of gamma ray signal should also become higher,
140 which means that we should set the threshold higher and the optimum length may become shorter than 6 cm.



142 FIG. 5. Pulse height distributions measured at FNL with changing the neutron energy generated through DD reactions where
143 the counts have been normalized by the measured counts with the ³He proportional detector.



145 FIG. 6. Pulse height distributions measured under irradiation of gamma rays from a ^{60}Co gamma ray source. The gamma
 146 rays were irradiated parallel to the axis of Sci-Fi.

147
 148 **B. Attenuation of scintillation photons through their transmission**

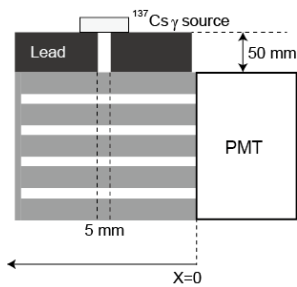
149 The scintillation photons can be scattered and absorbed in the transmission process inside the detector. If the Sci-Fi is long
 150 enough and acts as usual optical fiber, only the scintillation photons emitted to the direction over its critical angle can transmit
 151 to the PMT. However, as the Sci-Fi is short, the photons generated to all direction are reflected at the surface of Sci-Fi halls
 152 made through aluminum alloy and can transmit to the PMT. Through the transmission process, scattering and absorption at
 153 the interface between Sci-Fi and aluminum can affect the photon numbers to reach the PMT. To see the effect, we irradiated
 154 the gamma rays from ^{137}Cs gamma source with the configuration shown in Fig. 7, where gamma rays are collimated with a
 155 lead brick and irradiated perpendicularly to the detector with changing the irradiation position. The change of measured pulse
 156 height distributions are shown in Fig. 8 for the incident position of 0.25 cm, 1.25 cm, and 2.25 cm from the PMT from the
 157 PMT, where two types of the end reflectors were used, namely highly reflective type or matte black coating type. In Fig. 9,
 158 the counts over 100 ch are shown for the 3 incident positions. Each measured count with the ^{137}Cs gamma ray source at each
 159 position was normalized by the energy depositions simulated by using PHITS with the same configuration. From Fig. 8, the
 160 difference between the distributions of 0.25 cm and 1.25 cm or 1.75 cm is larger with a matt reflector compared to the case

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161 with the mirror reflector. From the results, we can conclude that the number of scintillation photons to reach the PMT is
162 affected also by the interaction position of neutron and the Sci-Fi. Especially when the end reflector is coated with matte
163 material, the difference is large. This should be the cause of broadening of pulse height distribution. However, the detection
164 efficiency changed gradually with the incident position, it seems that we should not necessarily pay excessive attention on to
165 the large signals from the nearest interaction which occurs in the nearest region of the PMT.

166 From the results shown above, we can conclude that there are two factors affecting the pulse height distribution measured
167 by the Sci-Fi detector, namely (a) self-shielding of 14 MeV neutrons by Sci-Fi itself and (b) attenuation of scintillation photons
168 during their propagation to the PMT. By the self-shielding effect, production of recoil protons near the PMT becomes fewer
169 for the 9 cm-long Sci-Fi than the 6 cm-long one. On the other hand, because of the difference of their length, more recoil
170 protons are generated in the 9 cm-long Sci-Fi at the positions far from the PMT. However, the photons generated at the
171 positions far from the PMT decrease before arriving the PMT, the resulting pulse height for that signal becomes lower and is
172 count in the pulse height region lower than the threshold. These effects totally result in the comparable counts for the 9 cm
173 and 6 cm Sci-Fi.

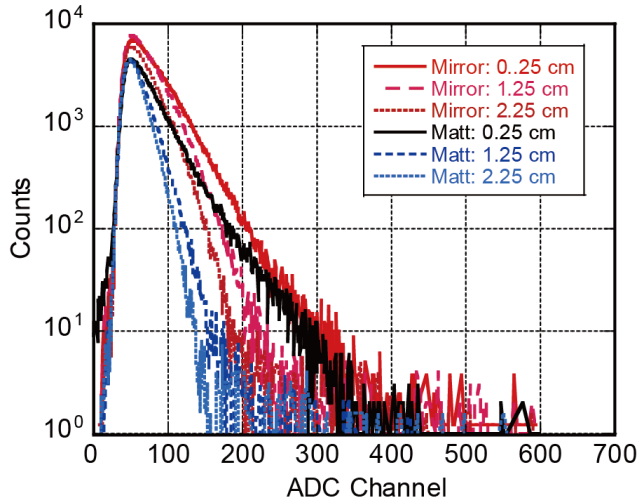
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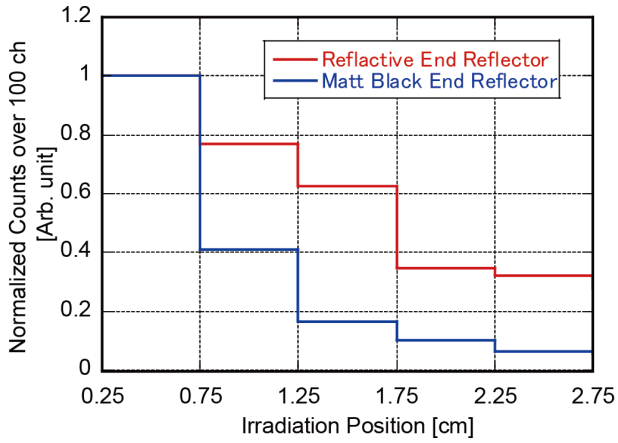
176 FIG. 7. Experimental setup to estimate the effect of interaction position on the pulse height distribution. The incident

177 position of ^{137}Cs gamma rays were changed.

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179 FIG. 8. Comparison of the pulse height distributions where the gamma ray incident position was set to 0.25 cm, 0.75 cm, and
 180 2.25 cm from the PMT. The two types of end reflectors were used, namely highly reflective type or matte black coating type.



182 FIG. 9. Relation between the incident position of ^{137}Cs gamma rays and the measured counts over the threshold of 100 ch. in
 183 Fig. 8.

184

185 **C. Comparison of measured and simulated detector responses**

186 To understand the phenomena occurring in Sci-fi detector, energy deposition in the detectors were calculated with PHITS.

187 In Fig. 10, the results are shown. As the Sci-Fi is very thin in diameter, once the incident neutron is scattered, it will go outside

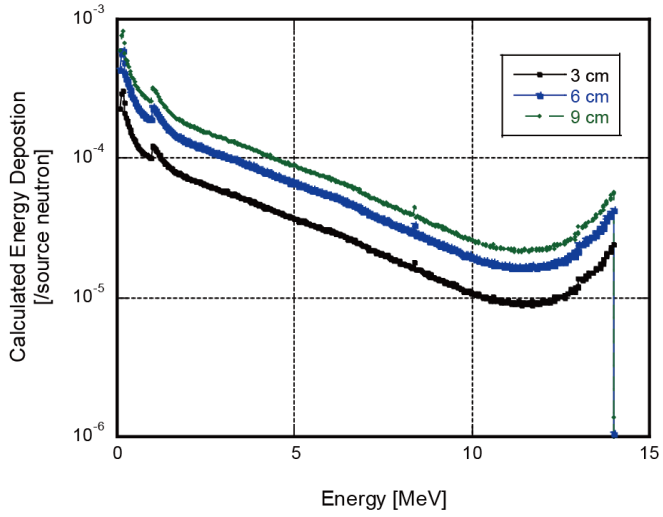
188 of the Sci-Fi. So, as the probability of secondary scattering is low, the shape of the pulse height distribution becomes
189 independent of its length and shows just a vertical shift for different length. Also, we can observe the steep edge at the highest
190 energy region for the detectors with three lengths. In the measured results, however, “knee” like distributions are observed at
191 that region. To reproduce the measured distribution by simulation, we considered the two phenomena as (a) production
192 efficiency of scintillation photons through interaction with protons compared with electrons and (b) probabilistic distribution
193 of measured scintillation photon numbers for each energy deposition. For (a), we calculated the light output from the Sci-Fi
194 by using the eq. (1) shown below which has been obtained by least square fitting of the data measured for plastic scintillator
195 (Saint-Gobain, BC-408)^[15].

$$L [\text{MeVee}] = 0.211 E_{\text{proton}}^{1.342} \quad (1)$$

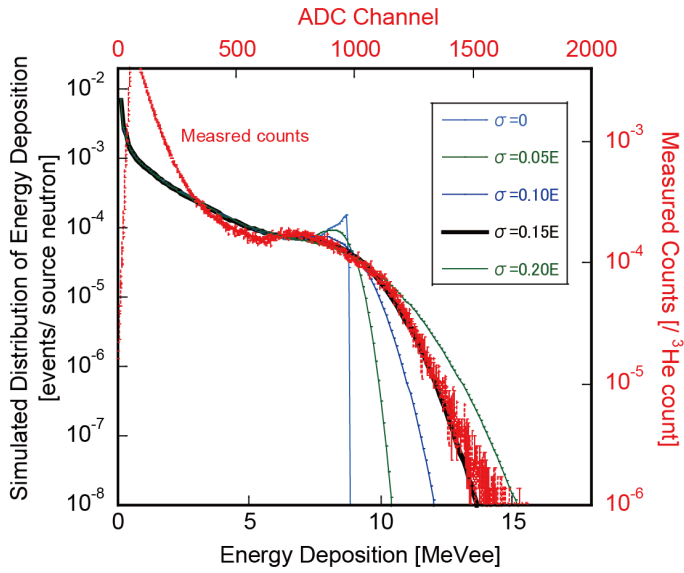
197 In considering the response of scintillation detector, namely the phenomenon (b), we assumed the Gaussian distribution
198 for the statistical distribution of the measured counts. In this assumption, the standard deviation: σ is proportional to the square
199 root of measured counts: N . If we assume that N is proportional to the electron equivalent energy deposition: E , σ can be
200 written as follows:

$$\sigma = K\sqrt{N} = A\sqrt{E}. \quad (2)$$

203 Based on the above relationship, we accounted the energy resolution in the simulated energy deposition and compared with the
204 measured values. Fig. 11 shows the results with assuming the proportional constant: A as 0, 0.05, 0.1, 0.15, and 0.2. From the
205 results, if we assume A to be 0.15, the simulated distribution shows the similar tendency with the measured one in the energy
206 range of above 5 MeV. However, in the lower energy region, the tendency is not the same. The measured pulse height is
207 lowered by the two phenomena as (a) low energy deposition because the proton was scattered to the direction with large angle
208 to the Sci-Fi axis and (b) loss of scintillation photons during their transmission process to the PMT. As, in the current simulation,
209 the phenomenon (b) is not considered, the simulated result does not show the same tendency with the measured one. To
210 estimate the response of Sci-Fi based system precisely, the transmission process of the scintillation photons should be
211 considered in the simulation.



213 FIG. 10. Simulated energy deposition in each Sci-Fi detector with three different lengths. The simulation was carried out by
 214 using PHITS code.



216 FIG. 11. Comparison of measured pulse height distribution and simulated electron equivalent energy deposition by PHITS
 217 for the 6 cm-long detector.

218

219 **V. SUMMARY**

220 To understand and enhance the properties of Sci-Fi detector for triton burnup experiments in magnetically confined fusion
221 plasmas, fundamental experiments and simulations have been carried out. From the results, it has been shown that the reflection
222 condition at the hole surface where the Sci-Fi are set is important. Also, it can be concluded that it is not necessary to make
223 Sci-Fi detectors longer than 6 cm, as the 14 MeV neutrons are shielded by the Sci-Fi itself and the energy deposition at far
224 positions from the PMT result in low pulse height due to scattering and absorption of scintillation photons in the transmission
225 process to the PMT. The information will be beneficial to design Sci-Fi detectors to be used at the fusion experimental devices.
226

227 **ACKNOWLEDGMENTS**

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229 (NIFS17KLEH068, NIFS15K0AH033) and the Japan-Korea Fusion Collaboration Program. The authors are also pleased to
230 acknowledge for the assistance of the FNL team and OKTAVIAN team in neutron measurements. ボディ文章中にFNL
231 も出てきますのでFNLも併記されたいかがでしょうか。
232

233 **REFERENCES**

- 234 ¹W. W. Heidbrink and G. J. Sadler, Nucl. Fusion 34, 535 (1994).
- 235 ²W. C. Sailor, Cris W. Barnes, R. E. Chrien, and G. A. Wurden, Rev. Sci. Instrum. 66, 898 (1995).
- 236 ³G. A. Wurden, R. E. Chrien, Cris W. Barnes, and W. C. Sailor, Rev. Sci. Instrum. 66, 901 (1995).
- 237 ⁴T. Nishitani, M Hoek, H Harano, M Isobe, K Tobita, Y Kusama, G A Wurdend, and R E Chrien, Plasma Phys.
238 Control. Fusion 38, 355 (1996).
- 239 ⁵H. Harano, JAERI-Research 97-060(1997) (in Japanese).
- 240 ⁶K. Ogawa M. Isobe, E. Takada, Y. Uchida, K. Ochiai, H. Tomita, A. Uritani, T. Kobuchi, and Y. Takeiri, Rev. Sci.
241 Instrum. 85, 11E110 (2014).
- 242 ⁷K. Ogawa, M. Isobe, T. Nishitani, S. Murakami, R. Seki, M. Nakata, E. Takada, H. Kawase, N. Pu, and LHD
243 Experiment Group, Nuclear Fusion, 58, 034002 (2018).

244 ⁸ K. Ogawa, M. Isobe, T. Nishitani, E. Takada, H. Kawase, T. Amitani, N. Pu, J. Jo, M. Cheon, J. Kim, M. Miwa, S.
245 Matsuyama and I. Murata, *Rev. Sci. Instrum.*, 89, 10I101-1 (2018).

246 ⁹M. Isobe , K. Ogawa, T. Nishitani , H. Miyake, T. Kobuchi, N. Pu, H. Kawase , E. Takada, T. Tanaka, S. Li, S.
247 Yoshihashi, A. Uritani, J. Jo, S. Murakami, M. Osakabe, and LHD Experiment Group, *IEEE Trans. Plasma Sci.*,
248 46, 6, 2050 (2018).

249 ¹⁰E. Takada, A. Fujisaki, N. Nakada, M. Isobe, K. Ogawa, T. Nishitani, H. Tomita, *Plasma and Fusion Research*,
250 11, 2405020 (2016).

251 ¹¹ I. Murata et al., *IAEA-TECDOC-1743* (Vienna: IAEA), 110–18 (2014).

252 ¹² M. Baba, M. Takada, T. Iwasaki, S. Matsuyama, T. Nakamura, H. Ohguchi, T. Nakao, T. Sanami, N. Hirakawa,
253 *Nucl. Instrum. Methods Phys. Res., A* 376, 115 (1996).

254 ¹³ T. Sato, Y. Iwamoto, S. Hashimoto, T. Ogawa, T. Furuta, S. Abe, T. Kai, P. Tsai, N. Masuda, H. Iwase, N.
255 Shigyo, L. Sihver, and K. Niita, *J. Nucl. Sci. Technol.* 55, 684 (2018).

256 ¹⁴ T. Nishitani, K. Ogawa, K. Nishimura, M. Isobe, *Plasma and Fusion Research*, 11, 2405057 (2016).

257 ¹⁵ J. Zhang, X. Ruan, L. Hou, X. Li, J. Bao, G. Zhang, H. Huang and C. Song, *Chinese Physics C*, 34, 988 (2010).

258