Identification of visible lines from multiply charged V8+ and V9+ ions

Identification of visible lines from Multiply charged W^{8+} and W^{9+} ions

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We present the investigation of new visible lines for W^{8+} and W^{9+} from the spectra observed using an electron beam ion trap. The analysis is based on the collisional-radiative (CR) modeling with fine-structure sublevels population kinetics. To ensure the identification of lines done by the CR model, we have also performed the accurate calculation of transition energies and transition probabilities within multi configurational Dirac-Fock using the GRASP2018. In the spectrum, most of the observed lines are assigned as magnetic dipole (M1) transitions belonging to the first two lower-lying configurations of these ions.

I. INTRODUCTION

Understanding the spectra from the lower charged states of tungsten is essential due to its application in several fields. In particular, tungsten is planned to be used as a plasma-facing component in the ITER divertor [1]. The temperature of the divertor plasma is expected to be few hundred eV, and in this temperature range, the most abundant tungsten charged states are W^{1+} to W^{15+} [2, 3]. Therefore, there will be a high possibility of observing UV and visible lines from these charge states from the ITER divertor region, which can be used in the divertor plasma diagnostics [4]. Especially, spectral studies of W^{7+} to W^{13+} ions are challenging as these ions have very complex atomic structures with the open $4f$ subshell, and competition of orbital energies between $4f, 5s$, and $5p$ electrons. For example, the one-electron Dirac-Fock energies of the $4f, 5s$, and $5p$ orbitals as a function of Z for Dy-like isoelectronic sequence (66 electrons) are calculated by GRASP2018 [5] and shown in Fig.1. It shows that around Z =74, i.e., W^{8+}, there is a considerable overlap of $5p$ and $4f$ orbital energies which results in the competition between $4f^{14}5s^25p^4$ and $4f^{13}5s^25p^5$ energy levels for the ground state. Due to such level crossing Berengut et al. [6, 7] have proposed forbidden transitions in W^{7+} and W^{8+} as prospective candidates for next-generation atomic clocks sensitive to variation in fine-structure constant. Not only on the application point of view, but spectral studies of these ions are also essential to test the limit of the existing theoretical methods.

To date, several studies are available on highly charged tungsten ions as the corresponding atomic systems (few electrons) are relatively simple. However, there is still a big void for lower charged tungsten ions from W^{7+} to W^{13+}, as pointed out by Ralchenko [8] and in data compiled by Kramida and Shirai [9]. In recent years, a couple of theoretical and experimental work is reported for W^{7+} ions. Berengut et al. [6] reported the energies for four levels from configuration interaction (CI) calculations, which matches the level sequence reported by Kramida and Shirai [9]. Ryabtsev et al. [10, 11] observed the spectrum from vacuum spark plasma and confirmed the ground state to be $4f^{13}5p^6$. Furthermore, the fine-structure splitting of the ground term $4f^{13}5p^6$ $2F_{5/2,7/2}$ was measured by direct observation of the M1 line with a compact electron beam ion trap (EBIT) in Tokyo (CoBIT) by Mita et al. [12]. The fine-structure splitting was also measured by Lu et al.[13]. Spectra of W^{11+}–W^{15+} in the 17–26nm region and a visible line from W^{11+} were measured from a compact EBIT in Shanghai (SH-HtscEBIT) [14]. In the visible range, 365-475nm few lines from charged state 8-28 were observed by Komatsu et al. [15] using CoBIT. However, for W^{8+} to W^{10+} ions, spectral data is still scarce, and no data available in the Atomic Spectra Database of the NIST [16]. Only for W^{8+}, sixteen lowest energy levels from $4f^{14}5s^25p^4$, $4f^{13}5s^25p^5$, and $4f^{12}5s^25p^6$ are theoretically calculated from FAC by Berengut et al.[6]. For W^{9+} and W^{10+}, no theoretical or experimental results are reported for energy levels or spectral lines.

To fill this void low charged tungsten ions (W^{6+}-W^{13+}) were produced and trapped in CoBIT [12, 17] previously. EUV spectra were recorded for W^{7+}-W^{13+} and visible spectra for W^{7+}-W^{9+}. However, in the visible range, only line from M1 transition between the fine-structure splitting of the ground configuration of W^{7+} was assigned while other lines presented in the spectra were unidentified. Therefore, in the present study, we analyze and identify the new lines from W^{8+} and W^{9+} ions with a detailed collisional-radiative (CR) model calculations. To aid the identification's reliability, we have also performed the structure calculation for these ions within the multi-configuration Dirac-Fock approach using GRSASP2018 [5].
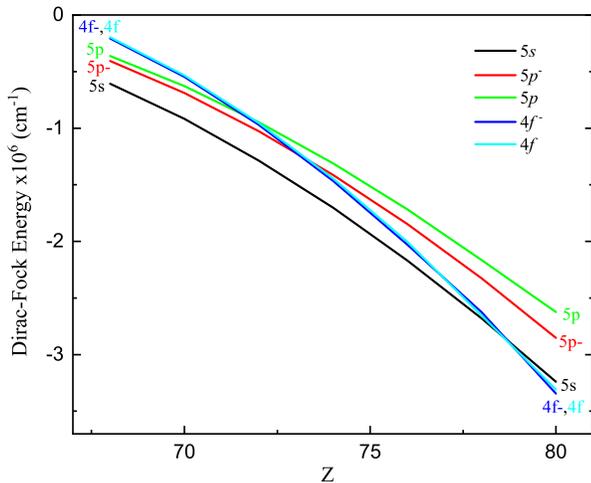
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FIG. 1. Dirac-Fock energies of Dy-like isoelectronic (N=66) sequence calculated by GRASP2018.

II. EXPERIMENT

The experiment was performed by using a compact EBIT (CoBIT) [18] at the University of Electro-Communications. For the details of the experimental setup, see our previous papers [12, 19]. Briefly, multiply charged tungsten ions were produced in CoBIT, consisting of an electron gun, a drift tube, an electron collector, and a high-critical-temperature superconducting magnet. The drift tube is composed of three successive cylindrical electrodes that act as an ion trap by applying a positive potential (typically 30 V) at both ends with respect to the middle electrode. The electron beam emitted from the electron gun is accelerated towards the drift tube while it is compressed by the axial magnetic field (typically 0.08 T) produced by the magnet surrounding the drift tube. The electron beam current was 3 to 5 mA depending on the electron beam energy. Tungsten was continuously introduced into the trap through a gas injector as a vapor of $W(CO)_6$. In order to maintain a good charge state distribution, the injection flow rate was carefully controlled by using a variable leak valve. Since there are no possible contaminating elements that are heavier than tungsten, ion dumping was not applied and, the tungsten ions were kept trapped during observations. The visible spectra were observed with a commercial Czerny-Turner type of spectrometer. A 300 mm^{-1} grating blazed at 500 nm was used for wide range observations, whereas a 1200 mm^{-1} grating blazed at 400 nm was used for wavelength determination. The data acquisition time was typically 20 min for wide range observations with the 300 mm^{-1} grating, and 40 min for high-resolution observations with the 1200 mm^{-1} grating. Wavelength calibration was done by using emission lines from several standard lamps placed outside CoBIT. The uncertainty in the wavelength calibration was esti-

mated from reproducibility to be about ±0.05 nm.

III. THEORETICAL CALCULATION

A collisional radiative (CR) model is developed for individual W^{8+} and W^{9+} ions using FAC 1.1.5 [20]. In the CR model, we assume that the various levels of the ions are interconnected through only collisional and radiative processes in the plasma. In the EBIT plasma, electron density is low, typically 10^{10-12} cm^{-3}, and in the present case, the electron beam energy is less than the ionization energy of W^{8+} and W^{9+}, on that note, we neglect various recombination processes such as three-body recombination and dielectronic recombination. In the present model, we consider the level population or depopulation by electron excitation, de-excitation, ionization, and radiative decay. In the steady state, for an excited level the rate balance equation with the normalization condition $\sum_j n_j = 1$, is given by

$$\sum_{\substack{i \\ i \neq j}} k_{ij} n_i n_e + \sum_{i>j} A_{ij} n_i - \sum_{\substack{i \\ i \neq j}} k_{ji} n_j n_e$$
$$- \sum_{i<j} A_{ji} n_j - n_j n_e k_{j+} = 0.$$

here, k_{ij}, A_{ij}, n_j and n_e are the electron impact excitation (de-excitation) rate coefficients, the transition probability from level $i \rightarrow j$, the population of the jth level and the electron density, respectively. Electron excitation rates are obtained by integrating the cross-section with the product of narrow Gaussian electron energy distribution function with 5eV FWHM as an EBIT has mono-energetic electrons. De-excitation rates are obtained by the principle of detail balance from excitation rates. In the present model, for W^{8+} total 4418 levels and ionization level generated from excitation from the reference $4f^{14}5p^4, 4f^{13}5p^5$ and $4f^{12}5p^6$ configurations to $5d, 5f$ and, for W^{9+}, total 14639 levels by excitation from $4f^{14}5p^3, 4f^{13}5p^4, 4f^{12}5p^5$ and $4f^{11}5p^6$ to $5d, 5f$ are considered as shown in Table I. Total more than a million cross-sections among all the states and millions of radiative decay channels (E1, E2, E3, M1, M2, and M3) for W^{8+} and W^{9+} are accounted in the model. All the atomic data for energy levels, transition probabilities, and cross-sections are calculated using the wavefunction obtained within the relativistic configurational interaction (RCI) method using FAC 1.1.5. The rate equations for all the considered levels are solved simultaneously to obtain the level populations and intensity ($I = E_{ij} * A_{ij} * n_j$). Spectra are obtained from line convolution with Gaussian functions having a standard deviation of 0.1 eV.

In FAC, the basis wavefunctions to calculate all the atomic properties are obtained from a single potential representing the screening of nuclear potential. The potential is optimized to the average electron clouds of configurations, which results in the less optimized potential

TABLE I. List of configuration considered in configuration interaction in FAC.

W8+	W9+
$5s^2 4f^{14} 5p^4$	$5s^2 4f^{14} 5p^3$
$5s^2 4f^{13} 5p^5$	$5s^2 4f^{13} 5p^4$
$5s^2 4f^{12} 5p^6$	$5s^2 4f^{12} 5p^5$
$5s^1 4f^{14} 5p^5$	$5s^2 4f^{11} 5p^6$
$5s^1 4f^{13} 5p^6$	$5s^1 4f^{14} 5p^4$
$5s^2 4f^{14} 5p^3 5(d,f)$	$5s^1 4f^{13} 5p^5$
$5s^2 4f^{13} 5p^4 5(d,f)$	$5s^1 4f^{12} 5p^6$
$5s^2 4f^{12} 5p^5 5(d,f)$	$5s^2 4f^{14} 5p^2 5(d,f)$
$5s^2 4f^{11} 5p^6 5(d,f)$	$5s^2 4f^{13} 5p^3 5(d,f)$
$5s^1 4f^{14} 5p^4 5(d,f)$	$5s^2 4f^{12} 5p^4 5(d,f)$
$5s^1 4f^{13} 5p^5 5(d,f)$	$5s^2 4f^{11} 5p^5 5(d,f)$
$5s^1 4f^{12} 5p^6 5(d,f)$	$5s^2 4f^{10} 5p^6 5(d,f)$
	$5s^1 4f^{14} 5p^3 5(d,f)$
	$5s^1 4f^{13} 5p^4 5(d,f)$
	$5s^1 4f^{12} 5p^5 5(d,f)$

for individuals. The average energy of each configuration group with the individually optimized potential and under the potential taking into account all configuration groups are obtained, and the difference between the two is used to adjust the final energy levels [20]. Further improvement is made by recalculating the average energy of each configuration under the radial potential obtained from a mean configuration (usually ground state configuration) from the OptimizeRadial module. The diagonal elements of the Hamiltonian calculated in the Structure module are then adjusted by the difference of the two average energies for each configuration. In the present case, for W8+ and W9+, including the mean average configuration corrections using $4f^{14} 5p^4$, $4f^{13} 5p^5$, $4f^{12} 5p^6$ and $4f^{14} 5p^3$, $4f^{13} 5p^4$, $4f^{12} 5p^5$ configurations, respectively, improves the level energies and level ordering matches with the GRASP calculations. However, the inclusion of the mean average configuration corrections decreases the accuracy in level splitting between individual configuration levels. It is found that the level energies of few lower levels are extremely sensitive to optimized potential due to substantial overlapping of levels from reference configuration.

Moreover, in the present RCI calculation, we have included the correlation effects only via important configurational interaction (which gives the best match to the experiment) to manage the model size as it includes calculation of a vast amount of collisional cross-section and decay probability data. Further, relativistic corrections from Breit interaction in the zero-energy limit for the exchanged photon and higher-order QED effects such as self-energy and vacuum polarization are added in a subsequent (RCI) calculation.

To ensure line identification from the CR model, we have also performed the multi-configuration Dirac-Fock

(MCDF) calculations using GRASP2018 [5]. We have calculated the energy of the first 30 and 144 lowest level and transition probability for M1 transitions in W8+ and W9+, respectively. In MCDF approximation, the wave function for an atomic state is approximated by an atomic state function (ASF), which can be expressed as a linear combination of configuration state functions (CSFs), which have the same angular momentum J and parity P [21, 24].

$$\Psi(\Gamma P J) = \sum_{i=1}^{n} C_i \phi(\Gamma_i P J).$$

Here, ϕ_i are the CSFs given as anti-symmetric products of the one-electron Dirac-Fock orbitals (Slater determinant). $(\Gamma_i P J)$ represents all information required to uniquely define the CSFs such as orbital occupation numbers, coupling, etc. n denoted the number of CSFs included in the expansion, and C_i's are the mixing coefficients. To obtain the final ASF's, we start with wavefunction calculation for multireference using a self-consistent-field (SCF) procedure based on the Dirac–Coulomb Hamiltonian and various correlation effects by including more CSF's layer by layer using active set approach [22, 23]. Moreover, corrections from the Breit interaction, i.e., transverse photon interaction in the low-frequency limit and the higher-order QED modifications due to self-energy correction and vacuum polarization (within a screened-hydrogenic approximation) are added in RCI calculations [25, 26].

Present systems W8+ and W9+ are near neutral ions having 3 holes and 4 holes, respectively, with $4f$ and $5p$ open orbitals. For these ions getting the convergence for the spectroscopic orbitals in the SCF procedure is very difficult. No convergence is found using initial estimation as screened hydrogenic functions and Thomas-fermi model. Therefore, first, we have used $DBSR_{HF}$ program [27] to generate orbital wavefunctions and used these as GRASP input for initial guess in SCF. DBSR is a B-spline version of a general Dirac-Hartree-Fock program. Then for all level calculation of the levels from reference configurations (for W8+ $5s^2 4f^m 5p^n$, $m+n=18$ and for W9+ $5s^2 4f^m 5p^n$, $m+n=17$) we include valence-valence and partial core-valence correlation effects via single double excitation from core valence to active set space. In the present calculation, all electrons are divided into two parts, electrons in $5p$ and $4f$ orbitals are taken as valence electron and in other inner orbitals as core electrons. Correspondingly, the correlation is taken as the interaction between the valence electrons and valence electrons with core electrons. In the present calculation, convergence is obtained by the following approach; we add the active set space in multi-reference Dirac Fock (DF) calculation layer by layer via single double (SD) excitation from $5s, 4f$, and $5p$ with a restriction that only one electron can excite from 5s at a time. While optimizing the outer layer, all the inner layers are fixed. We expanded the active set as follows:

W^{8+}	CSF's
DF+$(5d, 5f, 5g)_{SD}$	141,657
DF+$(5d, 5f, 5g)_{SD} + (6s, 6p, 6d, 6f)_{SD}$	485,857
DF+$(5d, 5f, 5g)_{SD} + (6s, 6p, 6d, 6f)_{SD} + (7s, 7p, 7d, 7f)_S$	488,622

W^{9+}	CSF's
DF+$(5d, 5f, 5g)_{SD}$	630,254
DF+$(5d, 5f, 5g)_{SD} + (6s, 6p, 6d, 6f)_{SD}$	799959
DF+$(5d, 5f, 5g)_{SD} + (6s, 6p, 6d, 6f)_{SD} + (7s, 7p, 7d, 7f)_S$	811309

In the case of W^{9+}, CSF's expansion was too large, so we put restrictions that orbitals in the active space are always doubly excited. Here, we have treated valence-valence and core-valence correlation via including CSF's generated from single and double excitation only from $5s, 4f$, and $5p$ orbitals. Sometimes deep core correlation by excitation from inner orbitals may also give some accountable contribution [22, 23]. However, due to complex and large expansion sets (even taking excitation only from $5s, 4f$, and $5p$), we did not include excitation from the inner core in RCI because of computational restrictions. Nevertheless, this approach can provide the accuracy in transition energies for line identification in the visible range, as shown in Tables II and III.

IV.   RESULTS AND DISCUSSION

In Fig. 2, experiential spectra observed with a 300 mm^{-1} grating are shown at various electron beam energies from 90-130eV in the visible region, and the lines, presented at 130eV are assigned to transitions in W^{8+} and W^{9+}. Since an EBIT has monoenergetic electrons, only those ionic state for which electron energy is higher than the ionization energy should be presented in the plasma. However, due to the excitation from metastable states, ions can be generated at an electron beam energy lower than the ionization threshold.

In the present case, it can be understood as follows, ionization energy from W^{7+} is 141eV, and from the CR model calculation, various metastable levels are found in W^{7+} (from $4f^{13}5p^55d^1$,$4f^{12}5p^65d^1$ levels) at around 40-60eV, that explains the presence of W^{8+} even at 90eV. In W^{8+} also there are various highly populated metastable states from $4f^{14}5p^4$, $4f^{13}5p^5$ and $4f^{12}5p^6$ (total 80% of the population) at 3-26eV, and few from $4f^{13}5p^45d^1$ and $4f^{12}5p^55d^1$ at around 50-57eV populated only 1%, and the ionization energy of W^{8+} is 160eV, which shows the presence of W^{9+} at 130eV. However, for ionization energy for W^{9+} is 180eV and the population is mostly concentrated in metastable states from $4f^{14}5p^3$, $4f^{13}5p^4$, $4f^{12}5p^5$ and $4f^{11}5p^6$ configurations situated up to 28eV from the ground, therefore at 130eV the most probable lines are from W^{8+} and W^{9+} ions.

In Fig. 3, we have shown simulated spectra for W^{8+} and W^{9+} at electron energy 130eV and electron density
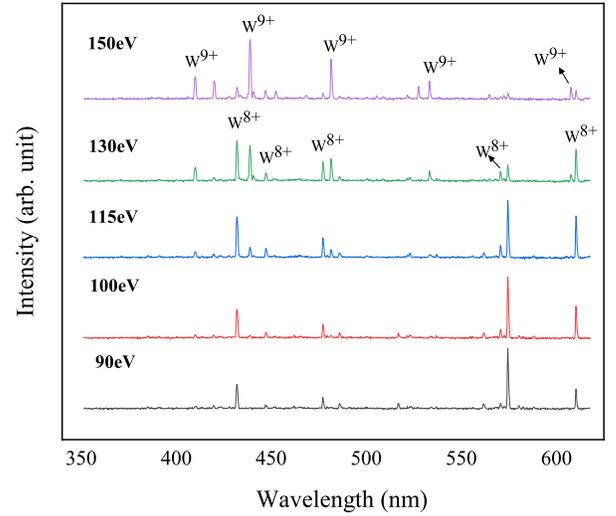

FIG. 2. Experimental spectra of tungsten ions in visible range obtained with a compact electron beam ion trap at various electron beam energies (written at the right side of the respective spectra).
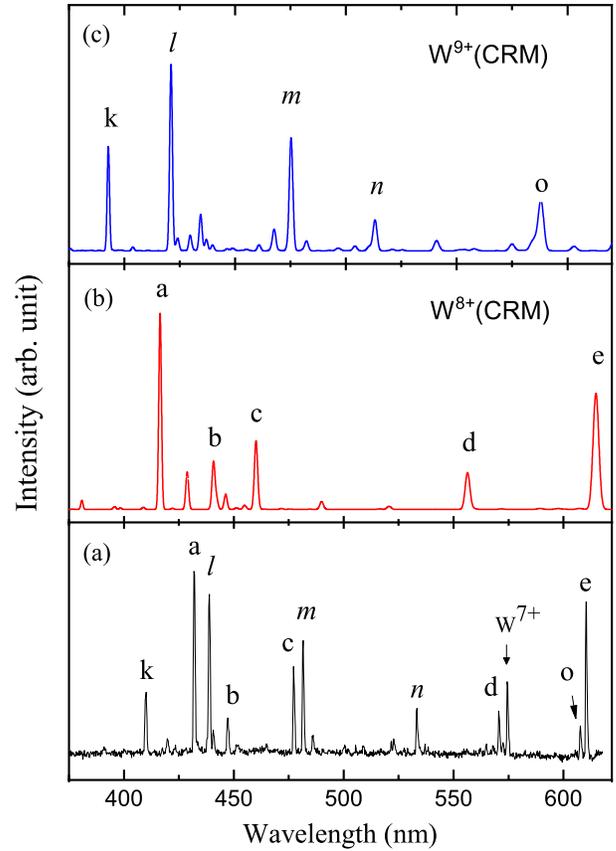

FIG. 3. Experimental (a) and simulated spectra of W^{8+} and W^{9+} ions (b) and (c), respectively, in the visible range at electron beam energy 130eV. The simulated spectra are obtained by CRM at electron density 10^{10} cm^{-3} using FAC.

TABLE II. Experimental and theoretical wavelengths (in nm) and theoretical transition probabilities A (s^{-1}) for transitions in Dy-like W^{8+}.

Label	Upper Level	Lower Level	λ_{exp}	λ_{th}				A		
				FAC	GRASP1	GRASP2	[6]	FAC	GRASP1	GRASP2
a	$((4\bar{f}^54f^8)_{5/2}(5\bar{p}^25p^3)_{3/2})_4$	$((4\bar{f}^64f^7)_{7/2}(5\bar{p}^25p^3)_{3/2})_5$	431.78	416.43	434.04	432.84	420.84	1.76e2	1.70e2	1.70e2
b	$((4\bar{f}^44f^8)_4(5\bar{p}^25p^4)_0)_4$	$((4\bar{f}^54f^7)_5(5\bar{p}^25p^4)_0)_5$	447.14	440.71	433.31	443.31	-	1.42e2	1.73e2	1.59e2
c	$((4\bar{f}^54f^8)_{5/2}(5\bar{p}^25p^3)_{3/2})_2$	$((4\bar{f}^64f^7)_{7/2}(5\bar{p}^25p^3)_{3/2})_3$	477.29	460.07	482.16	476.68	459.67	9.47e1	9.13e1	9.30e1
d	$((4\bar{f}^54f^7)_3)(5\bar{p}^25p^4)_0)_3$	$((4\bar{f}^64f^6)_4(5\bar{p}^25p^4)_0)_4$	570.52	556.32	544.50	571.25	546.24	8.22e1	9.06e1	8.45e1
e	$((4\bar{f}^54f^7)_5(5\bar{p}^25p^4)_0)_5$	$((4\bar{f}^64f^6)_6(5\bar{p}^25p^4)_0)_6$	611.17	614.72	604.75	619.74	596.24	1.05e2	1.18e2	1.10e2

TABLE III. Experimental and theoretical wavelengths (in nm) and theoretical transition probabilities A (s^{-1}) for transitions in Tb-like W^{9+}.

| Label	Upper Level	Lower Level	λ_{exp}	λ_{th}			A		
				FAC	GRASP1	GRASP2	FAC	GRASP1	GRASP2
k	$((4\bar{f}^54f^8)_{5/2}(5\bar{p}^25p^2)_2)_{7/2}$	$((4\bar{f}^64f^7)_{7/2}(5\bar{p}^25p^2)_2)_{9/2}$	409.66	398.65	427.35	408.19	1.17e2	1.38e2	1.30e2
l	$((4\bar{f}^54f^8)_{5/2}(5\bar{p}^25p^2)_2)_{9/2}$	$((4\bar{f}^64f^7)_{7/2}(5\bar{p}^25p^2)_2)_{11/2}$	438.68	424.41	450.71	442.33	1.68e2	1.85e2	1.78e2
m	$((4\bar{f}^54f^7)_5(5\bar{p}^25p^3)_{3/2})_{13/2}$	$((4\bar{f}^64f^6)_6(5\bar{p}^25p^3)_{3/2})_{15/2}$	481.55	478.78	477.74	483.71	2.25e2	2.21e2	2.17e2
n	$((4\bar{f}^54f^8)_{5/2}(5\bar{p}^25p^2)_0)_{5/2}$	$((4\bar{f}^64f^7)_{7/2}(5\bar{p}^25p^2)_0)_{7/2}$	533.20	514.88	555.06	536.70	1.16e2	1.30e2	1.25e2
o	$((4\bar{f}^54f^7)_5(5\bar{p}^25p^3)_{3/2})_{11/2}$	$((4\bar{f}^64f^6)_6(5\bar{p}^25p^3)_{3/2})_{13/2}$	608.41	596.08	590.63	604.53	1.16e2	1.13e2	1.11e2
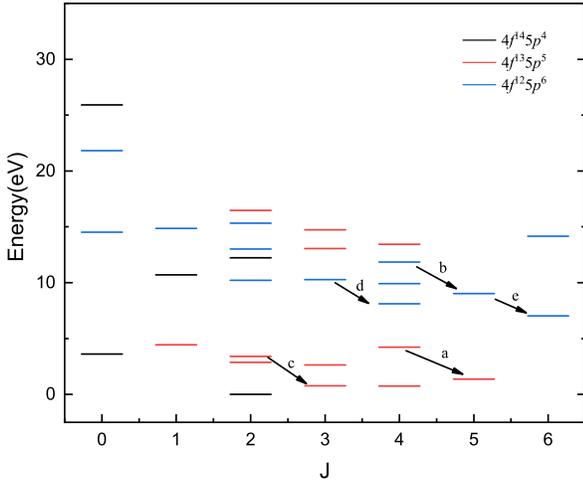
FIG. 4. Partial energy level diagram of W^{8+} with the lowest 30 energy levels calculated from GRASP2018. Different color show levels with different core. The vertical axis shows energy with respect to ground state $4f^{14}5p^4$ (J=2), and the horizontal axis shows total angular momentum. The arrow represents the M1 transition identified in Table II.
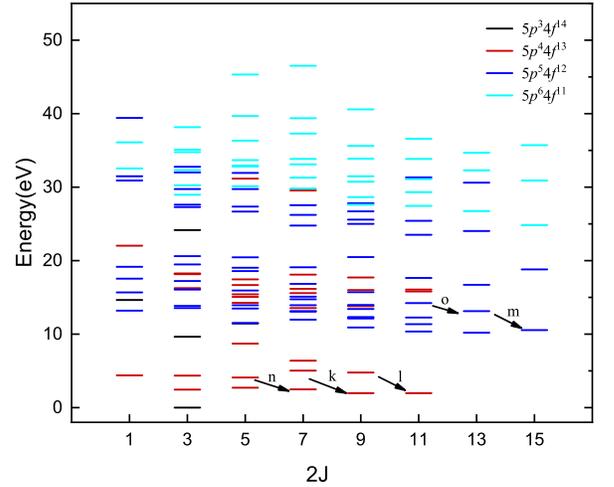

FIG. 5. Partial energy level diagram of W^{9+} with the lowest 144 energy levels calculated from GRASP2018. Different color show levels with different core. The vertical axis shows energy with respect to ground state $4f^{14}5p^3$ (J=3/2), and horizontal axis shows total angular momentum (2J). The arrow represents the M1 transition identified in Table III.

10^{10} cm^{-3} along with experimental spectra recorded at 130eV. It shows that spectral features from the model are similar to experimentally observed spectra, which allows us to identify the lines. From the CRM calculations, the observed lines are assigned as magnetic dipole (M1) transitions belonging to intra-transition between levels of $4f^{13}5p^5$ and $4f^{12}5p^6$ configurations in W^{8+} and $4f^{13}5p^4$

and $4f^{12}5p^5$ in W^{9+} ion.

The experimental wavelength of newly identified lines for W^{8+} and W^{9+}, along with theoretical values from FAC and GRASP, are presented in Tables II and III, respectively. Theoretically calculated transition probabilities are also given in the same tables. Levels are represented in $j - j$ coupled notation, where $n\bar{l}$ and nl
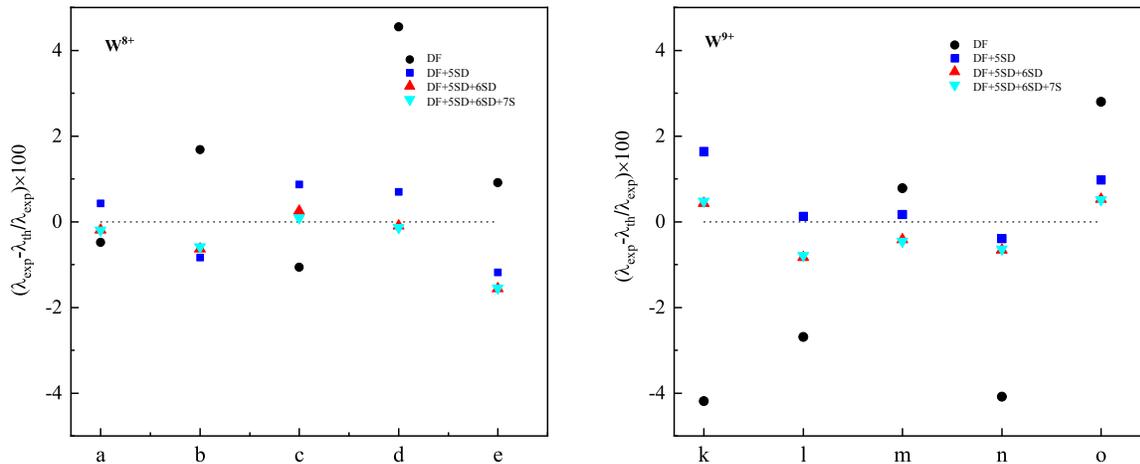
FIG. 6. Relative deviation in wavelengths obtained from the GRASP (for each expansion set) to the experimental wavelength for line labeled in Table II and Table III

corresponds to the shell with total angular momentum, $j = l - 1/2$ and $j = l + 1/2$, respectively. Lowest energy levels from reference configuration in W^{8+} and W^{9+} and identified transitions (labeled as a,b..) are presented in Figs. 4 and 5. Here, GRASP1 and GRASP2 represent DF+CI calculation only with reference configuration and including active space core-valence correlation corrections, respectively. For W^{8+} in Table II, we have compared the present wavelengths with the CI calculations of Berengut et al. [6]. It can be seen that theoretical values predicted from GRASP2 are more close to the measured wavelengths, and values reported [6] are similar to present FAC values. However, the difference between the transition energies reported in [6] and current theoretical values obtained from GRASP are within the 6000 1/cm uncertainty claimed in [6]. For W^{9+}, there are no previous experimental or theoretical data for transition wavelength or probability available for comparison with present values.

Table II shows that for W^{8+}, the theoretical wavelengths obtained from FAC are shifted in spectra with a maximum difference from the experimental value is 4%. Moreover, to ensure the identification, we have also compared FAC results with MCDF calculations from GRASP. From the difference between GRASP1 and GRASP2 wavelengths, it is clear that there is a significant contribution from the different correlations of the orbitals. The relative change in the transition wavelengths by adding active set expansion up to $(5d, 5f, 5g)_{SD}$ maximum 4% and further adding expansion up to $(6s, 6p, 6d, 6f)_{SD}$ and $(7s, 7p, 7d, 7f)_{SD}$ changes the wavelength by less than 1% and 0.15%, respectively. For W^{9+}, the theoretical wavelengths obtained from FAC are shifted in spectra with a maximum difference from the experimental value is 3.5%. The relative change in the transition wavelengths by adding active set expansion up to $(5d, 5f, 5g)_{SD}$ maximum 5.6% and further adding expansion up to $(6s, 6p, 6d, 6f)_{SD}$ and $(7s, 7p, 7d, 7f)_{SD}$ changes the wavelength by less

than 1.2% and 0.1%, respectively. From GRASP1 and GRASP2 results and mixing coefficients, it is noticed that for W^{9+} correlation contribution is larger than the W^{8+}. These results include the Breit and QCD correction contribution in low-frequency approximation, which increases the wavelengths approximately by 4%. The relative deviation in wavelengths obtained from the GRASP (for each expansion set) to the experimental wavelength is plotted in Figs. 6(a) and 6(b) for W^{8+} and W^{9+}, respectively. It shows that the expansion up to $7s, p, d, f$ orbitals reasonably explains the correlation. It can be improved by taking the correlation effect more accurately by including core-valence correlation from inner orbital and core-core correlation. Nevertheless, present accuracy provides the line identification very well.

In the present measurements, the identified lies are M1 transitions between states of the same configuration. They do not provide information about the relative position of the states of different configurations, which is essential to make any progress towards the new generation of atomic clocks. However, such measurements are very difficult since the transitions are very weak (strongly forbidden M1 transitions or E2 transitions). For example, the E2 transition from the ground state of W^{8+}, $4f^{14}5p^4 \ ^3P_2$, to $4f^{12}5p^6 \ ^3F_4$ with large sensitivity coefficient [6] has transition probability, almost two orders lower than for the currently observed lines. Therefore, at present, information about the relative position of the states of different configurations can be predicted from the theory only. In Table S1 and S2, we have tabulated level energies for lower lying levels of W^{8+} and W^{9+} calculated from GRASP2018 corresponding to Figs. 4 and 5 (see the supplementary tables)

V. SUPPLEMENTARY MATERIAL

However, it is hard to access the accuracy of transition energies in the absence of measurements.
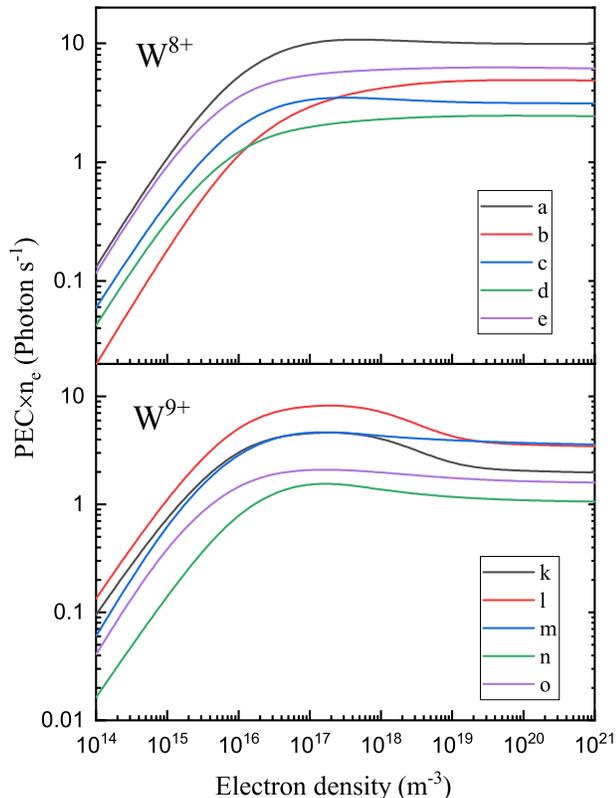
FIG. 7. Photon emission coefficient (PEC) as a function of n_e for the lines listed in Tables II and III estimated by CRM at Te=40ev for W^{8+} and Te=50eV for W^{9+}.

which W^{8+} and W^{9+} are most abundant) and shown in Fig. 7. In this case, the rate coefficients are calculated by using the Maxwellian electron energy distribution function. At higher electron density (10^{18} m^{-3}), PEC's are almost independent of electron density. We found that the present lines are as strong as the M1 lines of highly charged tungsten ions including W^{26+} observed in Large Helical Device (LHD) plasma at NIFS [28–31] for which PEC's are calculated to be the order of 10^{-18} m$^3/s$ at Te = 1keV and n_e = 10^{19} m^{-3}. Based on the visible emission line, they have demonstrated the evaluation of the ion temperature in LHD [30]. Therefore, such information for visible lines of W^{8+-9+} could be very important for the ITER plasma diagnostic purposes.

VI. SUMMARY

We identified new M1 lines from W^{8+} and W^{9+} ions in the visible range. To the best of our knowledge, these lines are measured for the first time. CR model is used to identify the lines, and also identification is ensured by comparing MCDF calculations from GRASP2018. The energy level calculations for the lower-lying levels belong to reference states in W^{8+} and W^{9+} are obtained by including core-valence correlations in a restricted manner. However, more rigorous calculations which account correlation effects completely can be further done as 2-hole system is studied in the recent theoretical evolution by Cheung et al. [32]. We hope that the present results are important in atomic structure calculations as well as very useful in fusion plasma diagnostics.

Furthermore, to see the importance of the identified M1 lines of W^{8+} and W^{9+} in the fusion plasma environment, we have calculated the photon emission coefficient (PEC) at various electron densities (10^{14}-10^{21} m^{-3}) and electron temperature 40 and 50 eV, respectively (at

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