

§21. Electrostatic Fluctuations in a Dimensionally Similar Low Temperature Plasma with Internal Transport Barrier

Takeuchi, M. (Dep. Energy Eng. Sci., Nagoya Univ.), Ikeda, R. (Dep. Energy Eng. Sci., Nagoya Univ.), Ito, T. (Dep. Energy Eng. Sci., Nagoya Univ.), Toi, K., Suzuki, C.

Electrostatic fluctuations have been measured by Langmuir probe from edge to core in low temperature plasma which is produced by 2.45 GHz microwaves at very low field less than 0.1T. It has the almost same dimensionless parameters, such as a normalized collisional frequency  $\nu_p^* \sim 0.05-1$ , plasma beta  $\beta \sim 0.001-0.02\%$  with those in high temperature one, except normalized ion gyro radius  $\rho_s^* \sim 0.015-0.04$ .

On the condition where the magnetic axis  $R_{ax}=97.4\text{cm}$ , toroidal magnetic field  $B_\phi=0.0613\text{T}$ , input heating power  $P_{\text{ext}}=17\text{kW}$  (oblique injection of O-mode) and hydrogen filling pressure  $p_H=7.0 \times 10^{-5}$  Torr, the transition in electron density  $n_e$  profile is often observed [1]. Time evolutions of  $n_e$  and line average density  $n_{e,2\text{mm}}$  measured by 2mm microwave interferometer are shown in Fig.1. During the decreasing phase of  $n_{e,2\text{mm}}$ ,  $n_e$  in the core region ( $\rho=0.3$ ) suddenly increases at  $t \sim 112\text{ms}$ , but that in the edge ( $\rho=0.7$ ) suddenly decreases. Radial profiles of  $n_e$ , electron temperature  $T_e$ , space potential  $V_s$  and the amplitude of normalized fluctuations before transition ( $t=95-105\text{ms}$ ) and after transition ( $t=120-130\text{ms}$ ) are compared in Fig.2. When the  $n_e$  profile is changed to a profile with a steep gradient inside  $\rho=0.7$ , the relative amplitude of density fluctuation  $\tilde{n}_e/n_e$  is obviously reduced at  $\rho=0.5-0.6$ . However,  $\tilde{T}_e/T_e$  and  $\tilde{V}_s/T_e$  are increased. Note that the profile of  $V_s$  changes from flat (radial electric field:  $E_r \sim 0$ ) to peak ( $E_r > 0$ ) in the core region. This clearly indicates the plasma goes into the electron root regime. Fluctuation driven particle flux  $\Gamma_{\text{turb}}$  and neoclassical particle flux  $\Gamma_e^{\text{nc}}$  before and after transition is shown in Fig.3.  $\Gamma_{\text{turb}}$  is much larger than  $\Gamma_e^{\text{nc}}$ .  $\Gamma_{\text{turb}}$  remains unchanged across the transition, but  $\Gamma_e^{\text{nc}}$  changes the sign across the transition. The flux change of  $\Gamma_e^{\text{nc}}$  is caused by the change to  $E_r > 0$ .

- [1] K. Toi et al., 13<sup>th</sup> Int. Toki Conf., Toki, 2003.
- [2] M. Yokoyama et al., J. Plasma Fusion Res. **79** (2003) 816

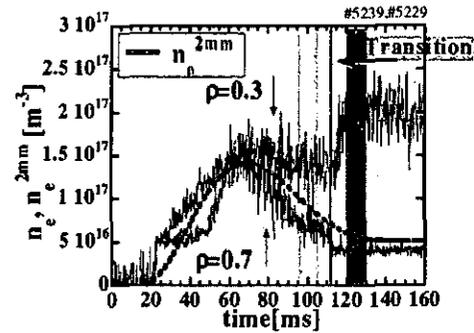


Fig.1 time evolution of electron density  $n_e$  and line average density  $n_{e,2\text{mm}}$

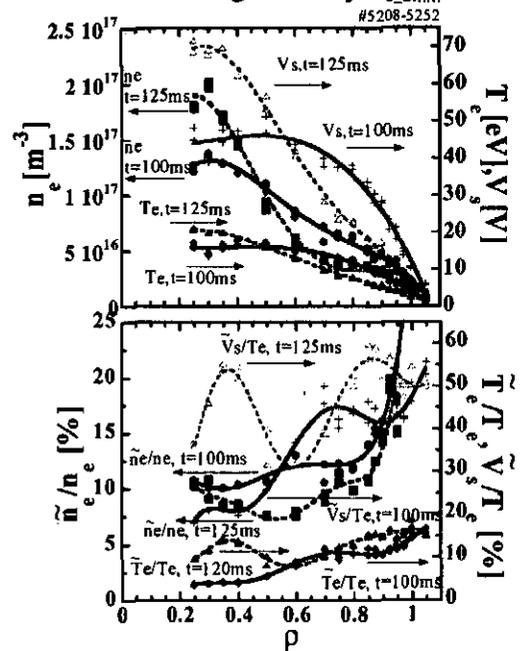


Fig.2 radial profiles of  $n_e$ , electron temperature  $T_e$ , space potential  $V_s$  and the amplitude of normalized fluctuation before transition ( $t=100\text{ms}$ ) and after transition ( $t=125\text{ms}$ )

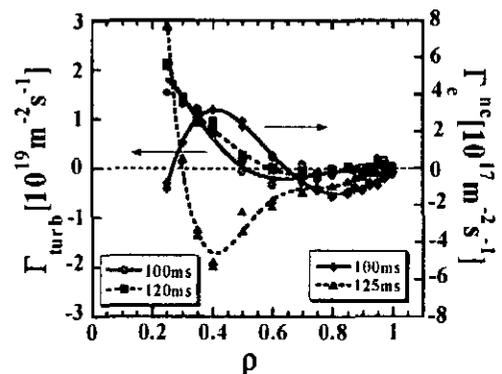


Fig.3 radial profiles of fluctuation driven electron flux  $\Gamma_{\text{turb}}$  and neoclassical electron flux  $\Gamma_e^{\text{nc}}$  before and after transition