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Bibliography of Electron and Photon Cross Sections with
Atoms and Molecules
Published in the 20th Century
– Carbon Dioxide –

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Bibliography of Electron and Photon Cross Sections

With Atoms and Molecules

Published in the 20th Century

—— Carbon Dioxide ——*

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(Gaseous Electronics Institute)

A bibliography of original and review reports of experiments or theories of electron and photon cross sections and also electron swarm data are presented for atomic or molecular species with specified targets. These works covered 17 atoms and 51 molecules. The present bibliography is only for carbon dioxide (CO₂). About 1,240 papers were compiled. A comprehensive author index is included. The bibliography covers the period 1901 through 2000 for CO₂. Finally, author's comments for CO₂ electron collision cross sections are given.

Keywords : CO₂ molecule, collision cross section, electron, elastic scattering, rotational, vibrational and electronic excitations, dissociation, ionization, photon, photoabsorption, photodissociation, photoexcitation, photoionization, electron swarm, drift velocity, diffusion coefficient, ionization coefficient, excitation and ionization energies, transition probabilities, lifetimes of excited states

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Introduction

History

This bibliography is the result of a continuing literature survey which was begun around 1970 and originally encompassed only electron collision cross section and electron swarm data. The organization responsible for continuing this survey is Nagoya Institute of Technology, Nagoya. From 1994, the work continued to Gaseous Electronics Insititute, Nagoya. In 1997, the collection of photon cross section references was begun. The search for references in both cases was retrospective and included all papers reporting measurements, theoretical calculations or reviews and data compilations of such cross sections and electron swarm data.

Scope

This bibliography contains references to original research papers which report experiments or theoretical calculations of cross sections for electron and photon collisions with carbon dioxide CO₂. The review papers on this subject are also included. Some nitrogen molecule cluster papers are included. Some conference reports, company or agency reports and PhD thesis are included. Carbon dioxide ion papers and positron collision papers are not included in principle.

Papers reporting the following data are included.

For electron collision cross section :

- 1) elastic scattering
- 2) rotational excitation
- 3) vibrational excitation
- 4) electronic excitation
- 5) dissociation
- 6) ionization
- 7) attachment
- 8) grand total scattering (sum of elastic and inelastic cross sections)
- 9) metastable nitrogen molecules
- 10) electron swarm parameters (drift velocity, diffusion coefficient)
- 11) excitation and ionization coefficients

For photon collision cross section :

- 1) photoabsorption
- 2) photoexcitation and fluorescence
- 3) photodissociation
- 4) photoionization

For some related data :

- 1) excitation and ionization energies
- 2) transition probabilities
- 3) lifetimes of excited states
- 4) others

The energy range for electron cross section data is usually 0 - 10 keV, but some higher electron energy papers are included. The wavelength range for photon cross section data is from microwave to X-ray. Most papers are concerned with infrared, visible and ultraviolet ray region.

The bibliography includes the papers published in the 20th century, from 1901 to 2000. Oldest paper in this list is given by J. S. Townsend (1901). So for this carbon dioxide bibliography, published papers from 1901 to 1999 are compiled by alphabetical order of the first author's surname of the paper. And the references published in 2000 and plus some old papers found very recently after compilation are added as "Addenda of References for Carbon Dioxide Molecule". In total, about 1240 papers are compiled in the carbon dioxide molecule bibliography.

Organization

This report consists of four parts : introduction, the bibliography and its addenda, author index, and some comments on electron collision cross sections.

Bibliography

In this section the complete citation for all references are given. At first following classifications are shown :

- E : Elastic collision
- R : Rotational excitation
- V : Vibrational excitation
- EX : electronic EXcitation
- D : Dissociation
- I : Ionization
- A : Attachment
- ME : MEtastable argon
- S : electron Swarm
- O : Others (photon cross sections and the others)

All authors' initials and surname, journal name, volume, inclusive pages and year of publication are given as well as the title, and some additional information in the square bracket []. E and T in the square bracket mean experiment and theory.

Bibliography for CO₂ are divided into two parts :

- Part 1. 1900 - 1999 p. 1 - 88
- Part 2. Addenda of References published in 2000, plus some old papers p. 89 - 106

Author Index

In this section all authors are listed alphabetically by surname. After each author's name is a list of page numbers indicating which references he or she authored or coauthored.

- Part 1. 1900 - 1999 p. 1 - 17
- Part 2. Addenda 1 p. 18 - 21

Some Comments on Electron Collision Cross Sections for CO₂

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References for CO₂ (1900 - 1999)

(Carbon dioxide, Carbonic acid)

E : Elastic collision.	R : Rotational excitation.
V : Vibrational excitation.	EX : Electronic excitation.
D : Dissociation.	I : Ionization.
A : Attachment.	QT : Grand total cross section.
S : Swarm.	α : Ionization coefficient.
O : The others.	[] : Additional informations.
	E : Exp., T : Theory.

The oldest paper in this list is given by J.S. Townsend (1901).

- A R. Abouaf, R. Paineau and F. Fiquet-Fayard : J. Phys. B9, 303-314 (1976) K
Dissociative attachment in NO₂ and CO₂.
[E, CO₂, NO₂; 3 - 4.8 eV for CO₂]
- O R. L. Abrams : Appl. Phys. Lett. 25, 609-611 (1974)
Broadening coefficients for the P(20) CO₂ laser transition.
[E, h ν , CO₂]
- O M. S. Abubakar and J. H. Shaw : Appl. Opt. 23, 3310-3315 (1984)
Analysis of carbon dioxide bands near 2.2 μ m. [E, h ν , CO₂]
- O M. S. Abubakar : PhD Dissertation, Ohio State University (1985)
Carbon dioxide band intensities in the 12.7 to 7.7 μ m region.
[E, h ν , CO₂]
- O M. S. Abubakar and J. H. Shaw : Appl. Opt. 25, 1196-1203 (1986)
Carbon dioxide band intensities and line widths in the 8 - 12- μ m region.
[E, h ν , CO₂]
- I B. Adamczyk, A. J. H. Boerboom and M. Lukasiewicz : Int. J. Mass Spectrom. Ion Phys. 9, 407-412 (1972) ·K
Partial ionization cross sections of carbon dioxide by electrons
(25 - 600 eV). [E, CO₂; CO₂⁺, CO⁺, O⁺, C⁺, O₂⁺]
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Mass spectrometric investigation of dissociative ionization of toxic
gases by electrons at 20 - 1000 eV.
[E, CO₂, N₂O, NO, CO; partial ioniz. c. s., 20 - 1000 eV]
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Emission cross sections of CO₂ by electron impact in the interval
1260 - 4500 A. II. [E, CO₂]

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Three-body electron attachment to a molecule. [review, CO₂, O₂, etc.]
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Electron transport coefficients in a weakly ionized plasma with
vibrationally excited molecules. [T, CO₂, N₂, CO]
- S N. L. Aleksandrov and I. V. Kochetov : J. Phys. D26, 387-392 (1993)
Electron transport parameters in a weakly ionized gas with vibrationally
excited molecules. [T, CO₂, N₂, CO; 1 - 300 Td]
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Electron rate coefficients in gases under non-uniform field and electron
density conditions. [T, CO₂, N₂, Ar]
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Ionization, attachment and negative ion reactions in carbon dioxide.
[E, CO₂]
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Molecular photoelectron spectroscopy. Part IV. The ionisation potentials
and configurations of carbon dioxide, carbon oxysulphide, carbon
disulphide, and nitrous oxide. [E, h ν , CO₂, COS, CS₂, N₂O]
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Resonant excitation of high vibrational levels by slow electron collisions.
[E, CO₂, H₂, N₂, CO]
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Study of triplet states and short-lived negative ions by means of
electron impact spectroscopy. [review, CO₂, N₂, CO, C₆H₆, etc.]
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Time of collection of electrons in ionization chambers. [, CO₂,
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Some measurements of electron drift velocities in compressed gases.
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Metastable fragment production following electron impact on CO₂.
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Observation of electron-neutral inverse bremsstrahlung in an electron-
beam-sustained discharge. [E, CO₂, N₂, H₂]
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On the cross section of low-energy electron collisions on CH₄ and CO₂.
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pathways involving ethene, propene and (E)-but-2-ene at atmospheric
pressure.
[E, CO₂; production of O(g) from CO₂, CO₂ + e → [CO₂]* → CO + O]
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Vibrational excitation of CO₂ by dipole interaction with slow electrons.
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Ionization of gases by 266-nm ultraviolet light.
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Zur Lumineszenz einiger Molekulgase bei Anregung durch schnelle
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Sur le seuil d'apparition des ions fragments produits avec excès
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[, hν, CO₂]

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Dissociative excitation of metastable fragments by electron impact on
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Evaluation of molecular quadrupole moments from broadening of microwave spectral lines. II. Calculation of the quadrupole moments and discussion.
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Author Index for CO₂ References

H. Abgrall 8
R. Abouaf 1
R. L. Abrams 1
M. S. Abubakar 1
Y. Achiba 41
M. Y. Adam 36, 66
B. Adamczyk 1
J. M. Ajello 1, 41
V. N. Akimov 73
P. Akther 40, 52
N. L. Aleksandrov 2
R. Alexandrescu 22
S. R. Alger 2
M. I. Al-Joboury 2
M. Allan 2, 22
G. Allicock 2
A. O. Allen 35
J. Allen 2
J. D. Allen 31
N. L. Allen 2
D. P. Almeida 2, 45
S. Alroy 3
S. H. Al-Shamma 65
H. Alvarez-Pol 3
M. A. Alves 26
S. V. Ambrosov 30
C. Anastasi 3
H. H. Andersen 63
T. Andersen 63
D. Andrick 3
Y. B. Anishchenko 3
H. Anton 3
Th. Antoni 3
J. Appell 3
G. Arena 4
C. B. Arends 22
G. S. Argyropoulos 3
E. Arie 3
N. Arimitsu 83
M. Armenante 4, 84
D. A. Armstrong 85
U. Asaf 4
L. Asbrink 27
R. K. Asundi 4
R. Atkinson 65
P. Ausloos 64
J. M. Austin 74
L. Avaldi 63
P. Ayotte 36
R. Azria 82
N. Azzi 87
T. Baer 4
P. S. Bagus 59
J. L. Bahr 4
V. A. Bailey 4, 5
J. K. Baird 15
A. D. Baker 83
C. Baker 83
C. J. Baker 5
J. H. Baker 5
A. K. Balan 30
K. K. Baldridge 70
K. L. Baluja 39
Y. B. Band 5
V. Berardi 4, 84
P. Barchewitz 24, 66
J. Bardelay 5
J. N. Bardsley 5
E. F. Barker 26
S. M. Barnett 5
D. M. Barrus 5
S. Barsotti 47
C. E. Bartky 12
R. Bartlett 66
M. J. Barton 5
G. Basavaraju 5, 51
A. D. Bass 36
H. Baumgartel 49
K. D. Bayes 6
M. Bayle 6
P. Bayle 6
E. C. Beaty 6, 28, 58, 59
K. Becker 20
U. E. Becker 82
K. Bederski 1
R. W. Bell 87
J. A. Benda 86

C. F. Bender 86
 W. S. Benedict 86
 D. C. Benner 6, 65
 S. W. Bennett 6
 V. Berardi 4, 84
 B. Th. Berendts 6
 L. -E. Berg 6
 J. Berkowitz 6, 23, 45
 H. J. Bernstein 57
 T. C. Betts 6
 K. D. Beyer 6
 R. A. Beyer 6
 M. S. Bhalla 7
 S. M. Bharathi 5, 51
 V. R. Bhardwaj 7
 K. G. Bhushan 5, 51
 S. E. Biagi 7
 N. K. Bibinov 7
 C. E. Bielschowsky 7, 19
 F. Biggs 65
 D. M. Binnie 7
 D. M. Bishop 7
 E. S. Bishop 7
 G. Black 74
 A. J. Blake 4
 R. L. Blake 5
 A. Blanchard 13
 G. S. Bloom 40
 C. E. Blount 7
 J. A. Boatz 70
 C. Bodere 8
 H. M. Boechat Roberty 7
 A. J. H. Boerboom 1, 10
 H. Boersch 7
 L. Boesten 43, 80
 A. Bogarski 7
 J. E. Boggs 56
 M. J. W. Boness 8
 V. Borodin 13
 W. L. Borst 56, 86
 T. E. Bortner 8
 L. F. Botelho 8
 P. Botschwina 75
 F. C. J. Bottiglioni 8
 L. Bouby 8
 R. Bouchard 39
 C. Boulet 17, 48, 53
 F. Bourbonneux 24, 66
 S. Bourquard 8
 A. Boyarski 8
 O. Boydston 30
 K. Boyer 50
 J. D. Bozek 8
 R. T. Brackmann 39
 N. E. Bradbury 9
 D. Bradley 9
 A. M. Bradshaw 70
 G. L. Braglia 9
 D. W. Branston 80
 M. Brennan 41
 D. C. Brenner 65
 L. M. Brescansin 37
 N. J. Bridge 9
 T. J. Bridges 9
 D. Briggs 8
 D. D. Briglia 9, 64
 C. E. Brion 9, 10, 14, 21, 26,
 28, 35, 48, 86
 R. B. Brode 10
 H. P. Broida 85
 J. P. Bromberg 10
 H. H. Brongersma 10
 I. K. Bronic 10
 H. L. Brooks 73
 H. L. Brose 10
 J. E. Brown 60
 L. R. Brown 10
 R. Browning 27
 E. Bruche 11
 P. J. Bruna 11
 C. R. Brundle 11, 83
 M. J. Brunger 11, 30, 40
 R. S. Brusa 78, 88
 R. Bruzzese 9
 W. A. Bryan 69
 A. D. Buckingham 9, 11
 S. J. Buckman 11, 30, 32, 82
 R. J. Buenker 11, 43, 61
 F. Bueso-Sanllehi 11
 M. O. Bulanin 11
 E. M. Bulewicz 11
 J. S. Bulger 11
 W. E. Bull 54
 B. R. Bulos 11
 V. P. Bulychev 11
 P. Bundgen 35
 D. E. Burch 11, 12
 P. R. Burchat 8
 A. J. Burek 5

I. Burghardt 12
D. J. Burns 52
P. D. Burrow 12
R. G. Buser 12
N. Bussieres 12
D. A. Byrum 65

I. Cadez 12
I. M. Cadez 12
J. Y. Cai 58
R. B. Cairns 12
G. E. Caledonia 12
R. Calilloni 30
R. Camilloni 31, 63, 76
M. W. P. Cann 13
P. Capezzuto 13
M. Capitelli 13
G. L. Caraffini 9
J. L. Cardner 4
J. -D. Carette 13
G. R. Carignan 73
R. W. Carlson 13, 48
T. A. Carlson 13, 31, 54
T. X. Carroll 13
A. I. Carswell 31
A. E. Carter 76
D. E. Carter 32
D. C. Cartwright 13
J. H. Carver 4
M. A. Casteel 3
A. W. Castleman 76
L. S. Cederbaum 21, 22
A. Cenian 13
R. Cesaro 4
J. E. Chaguri 58
G. Chambaud 75
K. C. Chambers 5
W. F. Chan 14
E. S. Chang 3, 14
T. Y. Chang 9
S. J. Chantrell 14
P. J. Chantry 14
R. E. Chapman 87
P. A. Chatterton 14
N. V. Cheburkin 30
A. Chedin 14
C. -F. Chen 14
C. L. Chen 61
C. T. Chen 50
F. Z. Chen 87
J. C. Y. Chen 14

Q. -M. Chen 87
S. -H. Chen 14
Z. Chen 26, 87
L. M. Cheng 7
A. Chernukho 13
J. Chiari 25
R. V. Chiflikyan 14, 23
B. K. Ching 16
W. H. Christiansen 3
B. M. Christopher 14
L. G. Christophorou 14, 15
A. Chutjian 32, 82
P. Cicman 15
R. Clampitt 15
S. A. C. Clark 53
D. C. Clary 34
C. R. Claydon 15
L. W. Cochsan 26
K. Codling 33
J. E. Collin 15, 36
L. A. Collins 15, 16, 55, 70
L. M. Colonna-Romano 24
N. Comanicu 22
J. Comer 16, 18
K. T. Compton 16
R. N. Compton 16, 43
A. Comunetti 16
V. J. Conti 16
D. C. Conway 16
L. W. Cochran 15
G. R. Cook 16, 85
C. D. Cooper 16
G. Cooper 14
M. A. Coplan 53
M. Coreno 63
C. Cornaggia 16, 17
S. J. B. Corrigan 17, 46
K. K. Corvin 17
C. Cossart-Magos 17, 75
V. W. Couling 17
C. Cousin 17, 48
C. Cousin-Lucasseau 17
J. D. Craggs 4, 7, 14, 17
F. Cramarossa 13
H. B. Crane 17
T. E. Cravens 18
B. L. Crawford 22
M. F. Crawford 76
R. W. Crompton 18, 30, 34, 37
A. Crowe 18, 89

F. Currell 18
 D. M. Curtis 18
 J. Cuthbert 18
 D. Cvejanovic 18
 S. Cvejanovic 18

 M. S. Dababneh 35, 46
 R. d'Agostino 13
 A. Dalgarno 18, 26
 N. Damany 18, 20
 N. Damany-Astoin 18
 C. J. Danby 22
 D. F. Dance 19
 N. P. Danilevskii 44
 A. M. Danishevskii 19
 A. Danjo 58
 D. Danner 3, 19
 A. Danno 38
 H. Date 33
 A. R. Davies 5
 D. K. Davies 19
 F. J. Davies 57
 F. J. Davis 84
 M. Davister 49
 G. Dawber 2
 E. F. Dawson 19
 C. J. Dedman 82
 F. J. de Heer 19
 J. Dehmer 50, 77
 J. L. Dehmer 19, 21, 37, 60, 66, 67, 73
 L. J. Dehmer 50
 P. M. Dehmer 19, 60
 E. P. de Lima 26
 G. Del Re 4
 M. J. DeLuca 19
 J. Delwiche 36, 66
 M. P. de Miranda 19
 W. B. De More 20
 A. J. F. den Boggende 20
 G. Denifl 15
 A. B. Denison 23
 D. M. Dennison 20
 W. Denzer 49
 L. De Reilhac 20
 E. Dershem 20
 L. Desesquelles 20
 S. Desi 57
 M. de Simone 63
 G. G. B. de Souza 7
 H. Deutsch 20, 51

 V. M. Devi 6, 20, 65
 A. D. Devir 20, 59
 M. Devoret 47
 Dibeler 20, 21
 D. M. Dickinson 7
 A. E. L. Dieperink 84
 P. Di Girolamo 84
 D. Dill 19, 21, 50, 77
 J. D. Dillard 21
 M. A. Dillon 47, 74
 M. Dillon 10
 R. L. Disch 11
 R. W. Ditchburn 21
 P. M. Dittman 21
 V. H. Dibeler 20, 21
 S. N. Dixit 53
 D. A. Dixon 82
 B. Dobson 43
 J. P. Doering 53
 A. I. Dolgin 51
 W. Domcke 21, 24
 F. H. Dorman 21
 R. Le Doucen 17, 48, 53
 S. R. Drayson 22
 H. Dreizler 31
 K. Dressler 34
 R. Dressler 22
 A. Dreuw 22, 63
 J. Drewko 78
 R. L. Dubs 53
 M. Dufay 20
 O. S. Duffendack 26
 A. Duguet 39
 D. Dumitras 22
 D. A. Dunmur 11
 I. Duran 3
 J. Durup 3
 J. Dutton 6, 22, 28
 P. P. D'yachenko 22
 A. Dymanus 6

 R. A. Eades 82
 D. E. Eastman 32
 D. F. Eggers 22
 H. Ehrhardt 3, 22, 44, 75
 M. Eidelsberg 17
 J. H. D. Eland 22, 23, 47
 J. H. S. Eland 18
 S. T. Elbert 70
 A. V. Eletskaa 23

M. T. Elford 11, 23, 32, 34
N. El Hakeem 52
R. E. Ellefson 23
R. E. Ellis 71
R. Ely 23
R. S. Eng 23
W. B. England 23, 24
W. England 23
P. Englander-Golden 9, 24, 64
W. N. English 24
W. C. Ermler 23, 24
D. D. Errett 24
J. Escada 26
F. Esposito 4
H. Estrada 24

I. I. Fabrikant 24, 41, 47
H. Falter 24
X. Fang 31
T. D. Fansler 24
J. F. Faris 54
C. B. Farmer 10
R. E. Farren 30
J. Farren 18
F. Farrenq 24
R. Farrenq 66
N. Feautrier 75
F. C. Fehsenfeld 25
J. Feldhaus 70
J. B. Fenn 48
J. Ferch 25, 77
E. E. Ferguson 25
R. Ferreira Marques 26
T. A. Ferrett 66, 67, 82
P. S. Ferry-Leeper 6
R. W. Fessenden 72
D. Field 14, 25
S. Filippi 4
S. V. Filseth 74
F. D. Findlay 13
M. Fink 25
F. Fiquet-Fayard 1, 8, 25, 30
O. B. Firsov 60
H. Fischle 26
I. M. Fishman 19
G. A. Fisk 34
W. L. Fite 39
S. H. Fleischman 26
G. D. Flesch 26

W. M. Flicker 46
G. W. Flynn 35, 58, 88
Ya. M. Fogel' 44
M. Fois 8
S. R. Foltyn 73
I. W. Fomunung 26
V. Y. Foo 26
D. W. Forester 15, 26
P. J. Fortune 23
M. C. Fowler 86
G. W. Fox 26
J. L. Fox 26
M. M. Fraga 26
J. L. A. Francey 26
J. L. Franklin 32, 46
G. W. Fraser 26
K. F. Freed 5
C. Freed 27
G. R. Freeman 39
J. P. Freeman 54
R. H. Freeman 27
L. C. G. Freitas 8
L. Frenkel 27
H. -J. Freund 27
R. S. Freund 27, 69
R. Frey 27
C. Fridh 27
A. A. Fridman 67
B. Fridovich 20
E. Friedlander 27
L. Frommhold 27
J. Fryar 27
R. Fuchs 28
J. H. Futrell 76

I. I. Galaktinov 28
J. W. Gallagher 28
I. Gallimberti 87
G. A. Gallup 28
T. E. Gangwar 35
K. S. Gant 15
D. Garand 65
A. V. Garchenko 30
G. Garcia 28
A. B. Gardner 85
J. L. Gardner 68
B. K. Garside 65
R. H. Garvey 31, 38
P. Gaspard 12
D. Gauyacq 28

J. Geiger 7, 28, 29
 E. P. Gentieu 29
 I. R. Gentle 29
 E. Gerjuoy 30
 E. V. Geroge 54
 E. T. Gerry 29
 A. G. Gershikov 29
 F. A. Gianturco 29, 81
 A. Giardini-Guidoni 30, 76
 J. C. Gibson 30, 82
 T. J. Gil 30
 A. L. Gilardini 30
 R. Gilpin 86
 F. M. Glaser 47
 V. M. Glazenzkov 30
 A. V. Glushkov 30
 W. A. Goddard III 86
 D. E. Golden 30
 T. I. Gombost 18
 J. M. Goodings 11
 A. Gopalan 47
 M. S. Gordon 70
 R. J. Gordon 88
 N. M. Gorshunov 30
 B. Gotchev 27
 S. Goursaud 30, 73
 M. C. Gower 31
 C. Graham 17
 P. Graham 31
 L. D. Gray 31
 A. E. S. Green 31, 38, 69, 77
 B. D. Green 12
 J. H. Green 31, 67
 M. A. Green 30
 F. R. Greening 31
 J. Gresser 70
 F. G. Gresteau 12
 F. Gresteau 12
 L. A. Gribov 31
 V. V. Grigor'yants 31
 F. A. Grimm 31
 J. Gripp 31
 R. Grisenti 78
 B. Grosswendt 84
 Z.-B. Group 44
 D. A. Gryvna 11, 12
 J. P. Gu 43
 W. Gudad 32
 A. G. Guidoni 31
 R. J. Gully 32
 S. K. Gupta 32
 M. V. Gur'ev 77
 T. Gustafsson 32
 P.-M. Guyon 36
 P. M. Guyon 4
 G. N. Haddad 23, 32
 A. Hadjiantoniou 14, 15
 O. F. Hagena 24
 R. D. Hake 32
 D. R. Hall 5
 R. I. Hall 12, 32
 A. Hamada 43, 68, 77, 79
 S. M. Hankin 31
 G. C. Hanna 24
 H. A. Hans 53
 J. E. Hardis 66, 67
 P. W. Harland 32, 83
 A. G. Harrison 32
 I. Harrison 32
 J. A. Harrison 69, 77
 P. Harteck 81
 F. Hartjes 33
 J. M. Hartman 33, 66
 K. O. Hartman 33
 J.-M. Hartmann 83
 J. M. Hartmann 33, 66
 C. N. Harward 33
 H. Hasegawa 33
 J. B. Hasted 8, 26, 33, 47, 88
 Y. Hatano 33, 51, 83
 P. A. Hatherly 33
 P. J. Hay 55
 M. Hayashi 33, 34
 M. A. Hayes 60, 72, 73
 G. N. Hays 34
 R. H. Healey 5, 34
 N. Hebel 44
 R. Hegerberg 34
 P. A. Heimann 82
 J. Heimerl 34
 J. Heintze 26
 F. Heitz 3
 B. Hemmerling 44
 W. Henkes 24
 H. J. Henning 34
 A. Henry 17, 48
 J. P. Henshaw 34
 G. H. Herzberg 34
 J. E. Hesser 34, 35

N. Hevel 75
 S. A. Hewitt 35, 88
 D. Hidalgo 51
 E. Hille 51
 G. Hirsch 43
 I. C. Hisatsune 33
 A. P. Hitchcock 35, 53
 K. R. Hoffman 35
 R. E. Hoffmeyer 35
 H. Hokazono 35
 M. L. Hoke 35
 D. M. P. Holland 72
 R. A. Holroyd 35
 W. Hoizer 57
 M. G. P. Homen 37
 A. Hopkirk 72
 D. G. Hopper 36
 M. Horani 36
 T. Horie 42
 M. Hoshino 68
 H. Hotop 47
 Y. -F. Hsieh 35, 46
 B. -L. Hu 55
 P. Huber 16
 M. J. Hubin Franskin 66
 M. -J. Hubin-Franskin 36
 R. D. Hudson 36
 M. A. Huels 36
 R. E. Huffman 36
 B. M. Hughes 36
 T. Hung 87
 W. T. Huntress 68
 G. S. Hurst 8, 14, 15, 37, 84
 H. E. Hurst 37
 N. S. Hush 42
 L. G. H. Huxley 37
 W. Hwang 37

 F. Iachello 37, 84
 M. C. Iasimone 4
 S. M. A. Ibrahim 9
 S. Ichimura 71
 I. Iga 37, 73
 T. Imamura 54
 M. G. Inghram 84
 E. C. Y. Inn 37, 85
 M. Inokuti 10, 19, 37, 38, 42, 60
 M. Inoue 38
 B. W. Irwin 50
 W. A. Isaacs 65

 I. Ishii 53
 T. Ishikawa 43, 79, 80
 I. Itikawa 43, 79
 Y. Itikawa 38, 43, 77, 78, 79, 80
 Y. Ito 81
 Yu. D. Ivanov 30
 T. Iwai 42, 83

 C. H. Jackman 31, 38
 W. M. Jackson 39
 F. M. Jacobsen 39
 A. Jaegle 39
 A. Jain 8, 29, 39
 D. K. Jain 39
 G. K. James 41
 G. Janzen 39
 A. Javan 27
 J. H. Jensen 70
 Y. Jiang 39, 77
 M. Jingchi 63
 H. W. Jochims 49
 T. L. John 39
 C. E. Johnson 40
 M. A. Johnson 19, 40
 H. L. Johnston 86
 W. H. Johnstone 40, 52
 W. M. Johnstone 40
 D. A. Jones 26
 E. G. Jones 32
 G. D. Jones 20
 K. D. Jordan 26, 66, 87
 K. N. Joshipura 40
 M. Joyeux 40
 H. Jucker 40
 D. L. Judge 13, 40, 41, 48, 62, 87
 K. Jung 3, 22, 44, 75
 M. Jungen 17
 J. Jureta 18
 A. S. Jursa 80

 S. Kadifachi 33, 88
 H. Kageshima 83
 H. Kallmann 27, 41
 O. F. Kalman 27
 K. Kameta 83
 N. Kameta 83
 S. Kaneko 68
 I. Kanik 41, 52
 A. Karawajczyk 6
 K. Karlsson 65
 G. P. Karwasz 88

G. Karwasz 78
 S. Katsumata 41
 W. E. Kauppila 35, 46
 M. Kawada 77
 A. K. Kazanskii 41
 A. K. Kazansky 41
 W. Kedzierski 41
 G. A. Keenan 19
 P. R. Keller 31
 H. G. Kerkhoff 82
 S. P. Khare 39, 42, 68
 S. Khare 51
 S. Ya. Khmel 42
 M. A. Khodorkovskii 51
 G. V. Khovrin 31
 E. B. Khudos 11
 M. H. Kibel 42
 L. J. Kieffer 42
 A. L. D. Kilcoyne 70
 D. A. L. Kilcoyne 42
 Y. -K. Kim 37
 K. Kimura 41
 M. Kimura 10, 43, 60, 68, 77,
 79, 80, 85
 Ma. Kimura 42
 Mi. Kimura 42
 A. D. King 43
 G. C. King 2, 32, 43, 72, 82
 G. W. King 31
 K. P. Kirby 43
 P. J. Kirkby 82
 J. Kirz 71, 87
 S. M. Kishko 43
 M. Kitajima 68, 77, 79, 85
 D. Kivelson 6, 81
 Y. Kiyama 58
 W. Klemperer 27
 L. E. Kline 60
 C. E. Klots 43
 K. N. Klump 44
 D. W. Knight 25
 G. Knoth 22
 P. H. Kobrin 82
 K. -H. Kochem 44, 75
 I. V. Kochetov 2, 44
 P. Kocian 8
 T. Koizumi 58
 Y. Kojima 58
 K. Kokubun 71
 R. Kollath 44, 63, 64
 T. Kondow 54, 83
 V. P. Konovalov 44
 M. D. Konstantinov 44
 N. Kontoleon 46, 84
 H. -M. Koppe 70
 V. T. Koppe 44
 C. Kosmidis 31
 H. Kossmann 27
 N. Kouchi 83
 A. P. Kouzov 44
 B. Kovac 44
 A. G. Koval' 44
 T. Z. Kowalski 44
 K. Kowari 60
 I. Koyano 45, 54
 D. N. Kozlov 44
 G. I. Kozlov 45
 J. U. Kozyra 18
 K. Kraus 45
 M. Krauss 45
 M. O. Krauss 31, 54
 A. Kresling 22
 E. Krishnakumar 45
 M. Krishnamurthy 45
 T. Kroin 45
 P. L. Kronebusch 45
 C. Kruger 57
 G. Kruppa 39
 K. Kuchitsu 54
 H. N. Kucukarpaci 45
 V. Kumar 4, 64
 A. Kuppermann 46
 S. E. Kupriyanov 46
 M. V. Kurepa 4
 M. Kurtz 18
 R. J. Kushlis 53
 M. S. Kushwaha 68
 K. Kutszegi 46
 C. E. Kuyatt 45
 M. Kuzumoto 46
 B. A. Kuzyakov 31
 Ch. K. Kwan 46
 P. Lablanquie 66
 P. Laborie 46
 N. LaCome 3
 A. Lahmam Bennani 46
 C. S. Lakshminarasimha 46, 73
 F. W. Lampe 46
 N. Lane 46

N. F. Lane 30, 46, 55, 86
 P. W. Langhoff 28, 60
 A. J. Langley 31, 69
 C. Larcher 28
 I. W. Larkin 8, 47
 U. T. Larnanna 29
 J. C. Larrabee 36
 W. Lasareff 27
 E. N. Lassetire 44, 47, 54, 72, 74
 F. Launay 17
 D. R. Laver 29
 G. M. Lawrence 47, 62
 S. Leach 17, 36, 40, 47
 E. Leber 47
 F. J. Le Blanc 80
 M. Lecas 46
 L. R. LeClair 47
 B. Leclerc 36
 S. Lederman 19
 K. W. D. Ledingham 31
 R. Le Doucen 17, 48, 53
 C. -H. Lee 48
 J. S. Lee 48
 L. C. Lee 41, 48, 62, 87
 N. Lee 48
 Y. T. Lee 85
 M. -T. Lee 8, 37, 45, 48
 H. Lehning 48
 S. W. Leifson 48
 K. Leja 67
 J. -L. Lemaire 87
 R. Lemus 37
 F. J. Leng 42
 D. A. Leonard 29
 S. -Y. Leu 14
 K. T. Leung 48
 R. Levi Di Leon 49
 N. E. Levine 49
 R. D. Levine 84
 B. R. Lewis 49
 J. Li 87
 J. C. Light 86
 A. A. Likal'ter 49
 J. W. Limbeek 49, 68
 F. Linder 49, 51
 E. Lindholm 27, 49
 D. W. Lindle 82
 B. G. Lindsay 77
 S. H. Linn 49
 Z. Y. Liu 63
 Li Xi 63
 A. N. Lobanov 49
 E. I. Lobodenko 61
 R. Loch 49
 P. Lodin 65
 M. Loewenstein 76
 R. K. Long 52
 R. Lorenzo 3
 J. C. Lorquet 50
 J. J. Lowke 50
 J. Lucas 45, 46, 49, 68, 84
 R. R. Lucchese 29, 50
 T. S. Luk 50
 P. Lukac 15
 M. Lukasiewicz 1
 S. Lunt 25
 D. L. Lynch 53
 M. G. Lynch 50
 M. Lynch 85
 O. M. Lyulin 61
 C. -M. Ma 50
 Y. Ma 50
 C. J. MacCallum 65
 M. A. MacDonald 72
 L. E. Machado 37
 K. Maciag 78
 S. R. Mackenzie 53
 R. G. A. R. Maclagen 83
 H. Mader 31
 B. Mahan 50
 S. Maji 5, 51
 S. V. Malinovskaya 30
 F. Manero 28
 A. W. Mantz 23
 C. Manzanares 51
 Z. -P. Mao 50
 D. Margreiter 51
 B. Marinkovic 78
 R. Maripuu 65
 T. D. Mark 15, 20, 51, 76, 86
 A. A. Markov 51
 P. Marmet 12
 G. Maroulis 51
 T. Masai 80
 C. Masche 25
 E. A. Mason 75
 N. J. Mason 5, 40, 52
 H. S. W. Massey 17
 A. Masuda 81

T. Masuoka 52, 68
 E. Mathieson 26, 52
 D. Mathur 7, 45
 F. M. Matsunaga 57
 J. L. Mauer 75
 D. P. May 2, 83
 H. F. Mayer 52
 J. M. Mayor 8
 S. F. Mazevet 30
 K. T. Mazon 37, 45
 D. J. McCaa 52
 T. McCanny 31
 D. C. McCollum 41, 52
 J. W. McConkey 2, 18, 41, 43, 47, 52
 J. H. McCoy 52
 T. K. McCubbin 23, 52
 K. E. McCulloh 53
 C. W. McCurdy 53, 65
 C. W. McCutchen 53
 R. McDiarmid 53
 J. M. McDonald 63
 S. P. McGlynn 63
 G. E. McGuire 13
 D. McKen 71
 V. McKoy 6, 48, 50, 53, 60, 78
 R. McLaren 53
 A. D. McLean 53
 S. M. McSweeney 72
 J. Mechlinska-Drewko 67
 D. Mehaffy 31
 G. Meigs 50
 G. G. Meisels 54
 R. Melissa 88
 S. Meloni 29
 V. Menoux 53
 J. E. Mentall 29, 53, 68
 F. Merkt 53
 R. H. Messner 53
 P. H. Metzger 16
 J. Meyer 4
 T. W. Meyer 53
 V. D. Meyer 47, 54
 S. E. Michelin 45
 K. Midorikawa 35
 S. R. Mielczarek 45
 L. D. Mikheev 7
 J. L. Miller 54
 P. Millie 66
 G. MilLOT 54
 F. Misaizu 54
 M. Misakian 54
 M. Mitchner 57
 R. K. Mitchum 54
 K. Mitsuke 54, 87
 A. K. Mnatsakanyan 54
 A. Kh. Mnatsakanyan 49
 W. E. Moddeman 54
 F. L. Mohler 20
 C. B. O. Mohr 54
 E. Molinari 13
 G. Mollenstedt 54
 J. E. Monahan 23
 K. Monahan 54
 A. Monfiles 34
 T. R. Mooney 52
 C. B. Moore 55
 L. A. Morgan 55
 W. L. Morgan 55
 S. Mori 77
 P. Morin 36, 66
 J. D. Morrison 21, 55
 M. A. Morrison 16, 30, 55
 A. L. Morse 40
 J. L. Moruzzi 56
 O. A. Mosher 46
 N. I. Moskalenko 56
 G. Mrotzek 25
 S. Mrozowski 56
 A. Z. Msezane 26
 D. Muigg 15
 F. Muller 25
 J. F. Mulligan 56
 R. S. Mulliken 56
 M. J. Mumma 54, 56, 73
 I. A. Munoz 51
 J. S. Murphy 56
 R. E. Murphy 12
 W. F. Murphy 57
 E. R. Murray 57
 C. S. Murthy 57
 K. Nagano 83
 A. F. Nagy 18
 G. P. Nagy 32
 L. Nagy 57
 T. Nagy 57
 M. S. Naidu 65, 73
 M. Nakamura 83
 Y. Nakamura 34, 57
 R. S. Nakata 57

H. Nakatsuji 57
 M. V. Narayan 71
 A. Ya. Nasarenko 29
 M. A. C. Nascimento 19
 P. Natalis 15
 V. G. Naumov 44
 D. R. Nelson 57
 I. Nenner 36, 66
 Yu. P. Neshchimenko 30
 D. Neumann 45
 W. R. Newell 5, 40, 69
 D. S. Newman 11
 A. S. Newton 15
 C. Y. Ng 26, 49
 B. Nguyen 46
 R. W. Nicholls 13
 J. Nickel 52
 J. C. Nickel 41
 F. H. Nicoll 54
 W. L. Nighan 57
 A. Nishimura 58
 H. Nishimura 57, 64, 68
 Y. Nishimura 58, 83
 B. Niu 19
 J. C. Nogueira 37, 58
 C. Nohre 65
 D. W. Norcross 60
 S. Nordholm 42
 D. Normand 16
 G. L. Nyberg 42
 K. J. Nygaard 73

 M. Obara 35
 A. A. Offenberger 14
 M. Ogawa 13, 16, 48, 58, 80, 85
 S. Ogawa 46
 Y. Oguma 58
 G. N. Ogurtsov 58
 Y. Ohmori 33
 H. Okabe 58
 M. Okamoto 68
 Y. Okamoto 59
 L. B. O'Kelly 37
 J. K. Olthoff 64
 K. Onda 58, 80, 81, 82
 J. A. O'Neill 58
 T. Ono 58
 C. B. Opal 58, 59
 V. P. Oppenheim 20
 U. P. Oppenheim 59
 O. J. Orient 59, 75

 S. V. Orlov 30
 V. E. Orlova 30
 Y. Oshima 59
 A. P. Osipov 59
 V. V. Osipov 44
 S. Oss 37, 78
 B. A. Osterlitz 23
 J. W. Otvos 59
 D. W. Overall 59
 J. Overend 60
 J. B. Ozenne 25

 P. M. Patel 40
 J. Pacansky 59
 J. L. Pack 59, 60
 N. T. Padiál 60
 N. Padiál 60
 A. Pagnamenta 60
 P. Pagsberg 3
 R. Paineau 1, 82
 P. Paoletti 29
 S. Pai 60
 L. A. Palkina 60
 L. Parenteau 36
 M. A. Pariseau 60
 J. H. Parker 50, 85
 D. A. Parkes 60
 J. E. Parkin 17
 A. C. Parr 60, 66, 67, 73
 G. R. Parr 60
 R. Parr 60
 M. Patapoff 20
 C. K. N. Patel 61
 P. M. Patel 40
 S. H. Patil 5
 R. R. Patty 12, 33
 W. B. Peatman 27
 J. Pebay 46
 A. Peluso 4
 J. Peresse 61
 V. I. Perevalov 61
 A. A. Perov 46
 M. Y. Perrin 33, 66
 K. I. Peterson 76
 W. K. Peterson 58, 59
 A. L. Petrov 61
 Z. L. Petrovic 61
 V. G. Pevgov 44
 V. G. Pevgov 61
 S. D. Peyerimhoff 61

S. D. Peyermhoff 11
 R. A. Phaneuf 80
 A. V. Phelps 11, 32, 50, 56, 59,
 60, 61, 62, 75
 E. Phillips 62, 87
 J. M. Phillips 23
 L. F. Phillips 6
 P. Pinson 5
 L. C. Pitchford 6, 28
 F. Piuze 62
 V. A. Pivovarov 62
 E. W. Plummer 32
 R. T. Poe 62
 V. Pol 35
 L. S. Polak 62
 H. M. Poland 62
 A. Policarpo 26
 H. Pollak 27
 L. Popova 44
 A. W. Potts 62
 M. C. Poulizac 20
 B. S. Prahallada 18
 A. M. Pravilov 7, 62
 M. R. Predtechenskii 84
 E. R. Preece 18
 A. L. Pregonzer 5
 W. Prepejchal 65
 W. M. Preston 62
 B. A. Prew 2
 D. A. Price 56
 W. C. Price 63
 K. C. Prince 63
 B. P. Pullen 54, 76

 T. Qinghua 63
 Y. Qui 63

 M. K. Raarup 63
 J. W. Rabalais 63
 M. Radle 22
 A. A. Radzig 63
 N. K. Rahman 29
 D. K. Rai 71
 W. Raith 25, 77
 D. Raj 42, 63
 F. A. Rajgara 7
 A. T. Rakhimov 59
 C. Ramsauer 63, 64
 J. Randell 25
 K. Randall 50

 M. V. V. S. Rao 64
 D. Rapp 9, 24, 64
 G. Rathenau 64
 P. Rawat 64
 F. H. Read 3, 43, 72, 82
 R. E. Rebbert 64
 T. Redish 18
 R. J. Rednall 53
 J. W. Reed 34
 J. A. Rees 2, 46, 64
 R. Reese 20
 R. R. Reeves 81
 D. F. Register 64, 82
 F. H. Reid 5
 J. Reid 65
 I. Reineck 65
 R. Reininger 4
 X. Ren 63
 T. N. Rescigno 65
 D. C. Reuter 73
 J. E. Reutt 85
 C. K. Rhodes 50, 53
 K. H. Richter 65
 E. K. Rideal 40
 B. A. Ridley 65
 F. F. Rieke 65
 W. Riemann 65
 M. E. Riley 65
 C. P. Rinsland 6, 10, 65
 A. V. Risbud 65
 J. S. Risley 83
 G. L. D. Ritchie 29
 D. Robert 17
 J. A. Roberts 27
 A. M. Robinson 65
 J. -M. Rocard 46
 C. Roche 54
 S. D. Rockwood 66
 V. I. Rodionov 30
 J. Romand 18
 L. Romano 9
 G. Romanowski 66
 P. L. Roney 13
 M. Roper 33
 B. Rosen 27, 34, 41
 B. J. Rosenberg 23
 L. Rosenmann 66, 71
 H. M. Rosenstock 53
 C. Rossetti 3, 24, 66
 P. Rosmus 75

K. J. Ross 47
 C. Rossetti 3, 24, 66
 A. R. Rossi 66
 B. Rossi 2
 F. Rostas 17, 87
 J. Rostas 28, 40, 87
 D. E. Rothe 52
 L. S. Rothman 66, 71
 L. S. Rothmann 71
 D. Roy 36, 66
 P. Roy 36, 66, 67
 W. Roznerski 9, 67, 89
 E. Rudberg 67
 J. B. Rudd 4, 67
 M. E. Rudd 37
 M. -W. Ruf 47
 J. R. Rumble 28
 V. D. Rusanov 67
 K. R. Ryan 31, 67
 V. A. Rykov 22

 E. H. Saayman 10
 H. T. Saelee 68
 T. Sagara 80
 N. Saito 8
 T. Sakae 68
 Y. Sakai 68
 S. Sakamoto 68
 Y. Sakamoto 68, 77, 85
 V. Saksena 68
 M. Salete Leite 26
 I. V. Samoilov 84
 J. A. R. Samson 12, 28, 52, 68, 85
 L. Sanche 12, 36, 69
 J. H. Sanderson 69
 M. G. Sanderson 3
 B. Sangi 69
 N. Sanna 29
 L. Sanson 18
 V. Santoro 4
 M. C. Sauer 69
 F. Sauli 71, 89
 T. Sawada 69
 M. N. Sawmy 69
 R. P. Saxon 37
 H. F. Schaefer 87
 G. C. Schatz 69
 M. Schein 20
 V. Scherr 63
 J. A. Schiavone 69
 R. N. Schindler 70

 J. Schirmer 21
 E. W. Schlag 27
 H. Schlumbohm 69, 70
 A. L. Schmeltekopf 25
 M. Schmidbauer 70
 B. Schmidt 26
 M. W. Schmidt 70
 M. Schmidt 16, 51
 V. Schmidt 27, 66, 67
 C. Schmiedekamp 25
 B. I. Schneider 15, 16, 70
 R. I. Schoen 70
 C. J. Schrijver 20
 U. Schucker 39
 E. Schultes 70
 G. Schultz 70
 G. J. Schulz 8, 16, 69, 70, 71, 75, 76
 R. Schumacher 70
 B. D. Schurin 71
 G. K. Schweitzer 54
 R. P. Schwenker 71
 D. Scutaru 71
 P. Seal 71
 T. J. Sears 71
 G. A. Segal 15
 H. J. J. Seguin 71
 P. Segur 87
 S. Sekine 71
 G. Senn 15
 A. M. Seregin 30
 L. Yu. Sergeeva 41
 L. Y. Sergeeva 41
 F. Sette 50
 T. K. Sham 71, 87
 R. G. Sharafutdinov 42
 V. N. Sharma 71
 A. Sharma 71
 T. E. Sharp 53
 W. E. Sharp 73
 R. L. Sharpless 74
 V. M. Shashkov 44
 J. H. Shaw 1, 35
 D. A. Shaw 72
 J. A. Sheehy 30
 D. E. Shemansky 72
 J. Shewchun 65
 J. C. Shiloff 47, 72
 N. Shima 83
 H. Shimada 81
 H. Shimamori 72
 I. Shimamura 72, 73

T. Shimanouchi 73
 H. Shimizu 71
 M. Shimosuma 33
 E. Shirakawa 58
 D. A. Shirley 82, 85
 G. V. Sholin 67
 O. B. Shpenik 88
 R. J. Shul 27
 I. O. Shul'pyakov 62
 T. W. Shyn 73
 M. C. Siddagangappa 73
 T. D. Sidorova 62
 K. Siegbahn 65
 J. Siegel 21, 50
 R. A. Sierra 73
 M. R. F. Siggel 60, 73
 A. Sillesen 3
 S. Silverman 86
 D. M. Simpson 63
 F. R. Simpson 52
 K. Singer 57
 Y. Singh 73
 R. P. Singhal 31
 J. M. Sirota 73
 V. N. Sivkov 73
 M. Sizum 30
 M. Sizun 73
 H. Sjogren 74
 J. D. Skalny 15
 A. M. Skerbele 47, 74
 M. F. Skinker 74
 V. V. Skubenich 43
 T. G. Slinger 74
 G. Sliwinski 13
 P. V. Slobodskaya 74
 D. I. Slovetsky 62
 J. H. Smart 35
 B. M. Smirnov 60, 63
 V. N. Smirnov 31
 A. L. S. Smith 74
 D. J. Smith 31
 K. A. Smith 77
 S. J. Smith 46
 H. D. Smyth 74
 A. V. Snegurskii 88
 A. V. Snegursky 88
 D. G. S. Snyder 20
 N. N. Sobolev 74, 75
 K. Soejima 58
 T. P. Softley 53
 W. Sohn 44, 75
 V. V. Sokovikov 74, 75
 T. Solovyeg 33
 A. J. Sommerer 73
 S. H. Southworth 66, 67
 U. Sowada 85
 D. Spence 70, 75
 F. E. Spencer 75
 A. Spielfiedel 75
 N. Spinelli 4, 84
 V. P. Spiridonov 29
 R. Spohr 75
 T. H. Spurling 75
 A. N. Srivastava 75
 S. K. Srivastava 59, 75
 W. Sroka 76
 W. Staib 39
 A. Stamatovic 15, 76
 S. Stangherlin 87
 M. Stankiewicz 33
 E. J. Stansbury 76
 A. Yu. Starikovskii 88
 W. L. Starr 76
 R. F. Stebbings 77
 G. Stefani 30, 31, 76
 T. S. Stein 35, 46
 I. T. Steinberger 4
 A. N. Stepanov 46
 K. Stephan 76
 D. P. Stevenson 59
 D. T. Stewart 85
 W. Stickei 7
 R. Stockbauer 84
 J. A. Stockdale 37
 J. A. D. Stockdale 15, 76
 T. Stoecklin 29
 E. J. Stone 56
 W. G. Stone 8
 A. R. Stopczynski 44
 G. Strakeljahn 77
 H. C. Straub 77
 D. J. Strickland 69, 77
 C. Stromholm 6
 E. C. G. Stuckelberg 74
 K. P. Subramanian 64
 A. F. Suchkov 49
 O. Sueoka 43, 77, 79
 H. Suhr 39
 J. J. Sullivan 12
 L. V. Sumin 77

H. Sun 77
 J. Sun 39, 77
 A. I. Suslov 44
 I. Suzuki 60, 77
 I. H. Suzuki 8
 S. Suzuki 54
 M. N. Swamy 77
 J. R. Swanson 77
 Cz. Szmytkowski 78, 89

 P. F. Taday 31, 69
 H. Tagashira 33, 68
 J. Taine 49, 66, 71
 T. Takahasi 58
 H. Takaki 43, 79
 S. Takasugi 58
 K. Takatsuka 50
 T. Takatsuka 78
 K. Takayanagi 78, 79
 M. Takekawa 43, 77, 79, 80
 K. H. Tan 9, 21
 H. Tanaka 43, 68, 77, 80, 85
 K. Tanaka 83
 M. Tanaka 46
 Y. Tanaka 36, 80
 S. M. Tarr 69
 H. Tashiro 35
 D. A. Tate 50
 H. E. Tatel 9
 R. Taubert 28
 H. Tawara 80
 S. S. Tayal 8, 39
 H. S. Taylor 15
 J. W. Taylor 31, 60
 J. L. Teffo 31
 J. -L. Teffo 61
 T. H. Teich 80
 D. Teillet-Billy 80
 J. B. Tellinghuisen 6
 A. Temkin 30
 J. Tennyson 80
 P. J. O. Teubner 30
 A. J. Thakkar 35, 51
 W. Thiel 80
 D. Thirumalai 80, 81
 R. V. Thomas 69
 T. D. Thomas 13
 B. A. Thompson 81
 D. G. Thompson 81

 J. P. Thomson 10
 A. M. Thorndike 81
 C. Tian 81
 R. Tice 81
 T. O. Tiernan 36
 J. Timmermans 33
 R. Tiribelli 30, 31, 76
 I. Tokue 81
 R. A. Toth 10, 81, 82
 J. S. Townsend 1, 82
 B. A. Tozer 17
 S. Trajmar 13, 32, 64, 82
 K. W. Trantham 30, 82
 W. Trela 66
 W. J. Trela 67
 M. Tronc 12, 82
 C. M. Truesdale 82
 D. G. Truhlar 58, 80, 81, 82
 J. S. Tse 71
 T. Tsuboi 83
 H. Tsuji 83
 M. Tsukada 83
 S. Tsuneyuki 83
 S. Tsurubuchi 83
 F. Tuffin 61
 J. Tulip 71
 D. W. Turner 2, 11, 83
 J. E. Turner 38
 D. C. Tyte 83
 P. Tzallas 31

 M. N. Uddin 46
 F. Udo 33
 K. Ueda 60
 M. Ukai 83
 M. A. Uman 49

 J. Vacquie 6
 C. Vallance 32, 83
 R. Vallauri 57
 R. J. Van Brunt 64
 P. J. M. van der Burgt 83
 J. A. Vanderhoff 6
 M. J. van der Wiel 35
 F. Vanoli 4
 O. S. van Roosmalen 84
 R. N. Varney 24
 H. Veenhuizen 65
 R. Velotta 4, 84

P. L. G. Ventzek 33
 H. F. A. Verhaart 84
 H. Versmold 57
 C. R. Vidal 81
 F. I. Vilesov 7
 D. Villarejo 84
 D. Vinciguerra 30, 76
 R. Vinciguerra 31
 A. S. Vinogradov 73
 I. P. Vinogradov 7
 L. E. Virr 84
 B. Vodar 18
 A. von Engel 5
 W. von Niessen 21
 E. Von Puttkamer 75
 R. E. Voshall 59, 60
 A. A. Vostrikov 84

 E. B. Wagner 37, 84
 A. C. Wahl 23
 V. Wahlgren 59
 E. Waibel 84
 N. Wainfan 84
 I. C. Walker 19
 J. A. Walker 21
 W. C. Walker 84
 A. D. Walsh 84
 L. Wan 39, 77
 K. P. Wanczek 66
 C. X. Wang 58
 L. -S. Wang 85
 J. M. Warman 85
 P. Warneck 85
 J. W. Warren 85
 R. W. Warren 85
 K. Watanabe 37, 57, 85
 S. Watanabe 68, 77, 79, 85
 T. Watanabe 38
 W. S. Watson 85
 R. B. Wattson 71
 T. S. Wauchop 45, 54, 85
 J. H. Weber 23
 A. B. Wedding 85
 H. V. Wedel 24
 E. Weigold 31
 G. L. Weissler 77, 84, 85
 J. Welch 21
 K. H. Welge 6, 45, 65, 86
 C. S. Weller 86
 E. J. Wells 20

 W. C. Wells 86, 88
 H. L. Welsh 76
 H. -J. Werner 75
 J. B. West 60, 73
 W. B. Westerveld 83
 R. C. Wetzel 27
 D. H. Whiffen 59
 B. L. Whitten 86
 J. L. Whitten 61
 W. J. Wiegand 57, 86
 L. Wiemann 25
 G. R. Wight 86
 P. G. Wilkinson 86
 A. W. Williams 5, 16
 D. Williams 11, 69
 I. D. Williams 69
 P. I. Williams 14
 T. A. Williams 62
 C. Winkler 86
 C. Winstead 48
 C. L. Winstead 30
 N. W. Winter 86
 B. H. Winters 86
 K. Wittmaack 28, 29
 L. Wojcik 1
 C. F. Wong 86
 W. C. Wong 6
 R. E. Wood 55
 J. M. Woolsey 52
 C. Y. R. Wu 87

 Q. -H. Xu 87

 S. Yagi 46
 J. Yan 50
 B. X. Yang 71, 87
 X. -F. Yang 87
 J. T. Yardley 55
 I. D. Yaroshetskii 19
 V. S. Yarunin 87
 D. Yeager 23
 D. C. S. Yee 9
 W. -K. Yen 14
 L. Yin 68
 H. Yoshida 87
 K. Yoshida 33
 Y. Yoshida 87
 C. Young 22, 87
 D. R. Young 87
 L. D. G. Young 66

M. Yousfi 87
M. S. Yurev 87

R. N. Zare 40
A. N. Zaviropulo 88
A. Yu. Zayats 46
A. Zecca 78, 88
M. Zelikoff 37, 85
M. E. Zhabotinskii 31
M. C. Zhang 63
J. Zhou 50
L. Zhu 35, 88
Y. -F. Zhu 88
A. I. Zhukov 88
J. P. Ziesel 25
J. -P. Ziesel 25
R. Zietz 76
T. M. Zimkina 73
U. Zimmermann 57
E. C. Zipf 56, 86, 88
N. Zonjee 33
M. Zubek 78, 88
A. P. Zuev 88

Author Index for CO₂ References. Addenda 1

- W. Addo-Asah 100
H. Agren 99
M. Alacid 90
C. J. Allan 91
M. Allan 89
M. Allen 91
W. D. Allen 91
D. A. Allison 91
G. Amat 91
C. Amiot 98
R. D. Amos 91
A. D. Anbar 91
A. Anderson 91
R. Anttila 99
P. Arcas 91, 98
P. Archirel 98
E. Arie 91, 98
R. Azria 101
- D. Bailly 91, 92
D. J. Baker 101
K. C. Baker 101
A. Baldacci 92
L. S. Ballomal 91
A. Barbe 100
E. F. Barker 92
S. Barsotti 90
R. J. Bartlett 95
A. Bastida 90, 100
M. R. Battaglia 91
C. E. H. Bawn 97
W. S. Benedict 92
D. C. Benner 92, 98, 100
J. Berkowitz 94
N. Berrah 90
R. Berson 92
A. Bastida 90, 100
T. G. Beuthe 92
J. H. Beynon 93
C. E. Bielschowsky 100
S. Bodeur 99
R. W. Boese 102
C. Bohne 101
B. Bonnet 96
J. D. Bozek 90
- M. Breitenstein 92
A. G. Brenton 97
K. G. Brown 99, 101
L. A. Brown 100
L. R. Brown 92, 94, 100
R. S. Brusa 89
D. E. Burch 92
- R. F. Calfee 92
R. Camilloni 97
C. Camy-Peyret 92, 100
M. Capitelli 93
S. Carter 93
L. S. Cederbaum 93, 102
V. Cermak 93
An-Ti Chai 91
J. -S. Chang 92
Yu. Chang 103
A. Chedin 93, 96
X. J. Chen 89
Z. Cihla 93
D. A. Clabo 91
I. Coleman 101
J. E. Collin 96
R. N. Compton 93, 95
R. G. Cooks 93
C. D. Cooper 93
A. G. Csaszar 91
J. Cuisenier 98
M. Cuisenier 91
- P. Dahoo 95
V. Dana 93, 95
L. E. Danilochkina 102
J. Delewiche 99
G. De Maria 94
T. E. Derviz 102
G. G. B. de Souza 100
V. M. Devi 98
M. Devoret 94
G. H. F. Diercksen 99, 102
W. Domcke 93, 94
H. D. Dowling 94
J. Dupre-Maquaire 94

D. P. Edwards 101
 J. H. D. Eland 94, 96, 98
 P. W. Erdman 94
 P. Erman 103
 M. P. Esplin 94, 104
 H. Estrada 94

I. I. Fabrikant 90
 R. Farrenq 91
 G. H. Fattahallah 95, 99
 D. Field 89
 M. Fink 98
 U. Fink 100
 J. -M. Flaud 100, 101
 P. Fournier 98
 D. C. Frost 95
 K. Furuya 89

R. R. Gamache 100
 U. Gelius 91, 99
 A. G. Gershikov 95
 F. A. Gianturco 89
 J. R. Gillis 100
 A. Goldman 100, 101
 H. R. Gordon 95
 L. D. Gray 95
 A. E. S. Green 95
 W. F. Grieder 101
 P. Gross 100
 N. E. Gruner 99
 D. A. Gryvnak 92
 G. Guelachvili 91, 93, 95, 99
 G. L. Gutsev 95

M. Hamdan 97
 A. Hamdouni 93, 95
 N. C. Handy 93
 J. Hansen 95
 A. Henry 95
 A. Herzenberg 95, 96
 A. P. Hitchcock 96
 M. L. Hoke 94
 M. Horani 96
 V. -M. Horneman 96, 97
 H. Hotop 90
 S. Hsieh 96
 M. J. Hubin-Franskin 99
 M. -J. Hubin-Franskin 96
 R. D. Hudson 96
 R. H. Hunt 94, 101

R. J. Huppi 94
 N. Husson 96, 100

W. A. Isaacs 90
 Y. Itikawa 89, 90
 A. Javan 98
 C. C. Jia 89
 J. A. Joens 96
 G. Johansson 91
 J. W. C. Johns 96
 D. Johnson 95
 K. Jolma 96, 97
 P. Jonathan 97
 N. C. Jones 89

A. Karawajczyk 103
 G. P. Karwasz 89
 J. Katihabroa 96
 W. E. Kauppila 100
 J. Kauppinen 96, 99
 J. K. Kauppinen 97
 A. K. Kazanskii 97
 V. A. Kazbanov 102
 M. J. Kelly 98
 X. Kezun 103
 M. Kimura 89, 90
 M. Kitajima 89, 90
 J. Klein 98
 U. Koble 103
 I. V. Kochetov 97
 T. Kondow 97
 Q. Kou 93
 W. P. Kraemer 102
 E. Kukk 90
 K. C. Kulander 97
 N. A. Kurnit 98
 C. K. Kwan 100

P. Lablanquie 98
 A. Lacis 95
 M. Lacy 97
 R. Lanquetin 92
 S. Leach 94, 96
 S. Lebedeff 95
 E. Leber 90
 R. E. Leckenby 97
 N. Leclercq 90
 P. Lee 95
 S. T. Lee 95
 T. J. Lee 98

P. Letardi 97
 J. Li 97
 E. Lindholm 97
 B. G. Lindsay 90
 Y. Liqiang 103
 W. Lochte-Holtgreven 97
 S. L. Lunt 89

 J. P. Maier 98
 J. P. Maillard 91, 98
 A. G. Maki 99
 L. Malaspina 94
 V. Malathy Devi 98, 100
 L. Malegat 101
 J. -Y. Mandin 93
 N. Martensson 99
 J. M. L. Martin 98
 S. T. Massie 100, 101
 D. Mathur 100
 A. Matsuo 89
 R. J. Mawhorter 92
 S. Mazevet 90
 J. J. McClelland 98
 T. K. McCubbin 95
 C. W. McCurdy 90
 C. A. McDowell 95
 H. Meyer 92
 H. -D. Meyer 90
 R. Middleton 98
 P. Millie 98
 C. Miron 90
 K. Mitsuke 97
 M. H. Mittleman 98
 M. Molinari 93
 J. -P. Monchalin 98
 L. A. Morgan 90
 M. Morillon-Chapey 93
 P. Morin 90, 99
 M. A. Morrison 90
 R. S. Mulliken 98
 J. R. Murray 98

 R. M. Nadile 101
 H. A. Nair 91
 I. Nenner 98
 R. K. Nesbet 90
 A. Nilsson 99
 D. Nordfors 99

 J. Oddershede 99

 T. Ogawa 89
 A. E. Orel 90
 G. Ouyang 89

 D. A. Parkes 99
 R. Paso 99
 R. R. Patty 92
 L. L. Peng 89
 A. Perrin 100, 101
 W. B. Person 99, 101
 A. A. Peshkov 90
 D. Peters 99
 F. R. Petersen 99
 V. Piacente 94
 A. S. Pine 99
 P. Pinson 94
 E. K. Plyler 92, 101
 A. W. Potts 95, 99
 J. Pritchard 101
 D. A. Przybyla 100

 E. Rachlew-Kaline 103
 F. A. Rajgara 100
 D. H. Rank 100
 G. Ravindra Kumar 100
 W. T. Rawlins 101
 R. B. Remington 91
 A. Requena 90, 100
 T. N. Rescigno 90
 C. K. Rhodes 98
 D. J. Richardson 100
 D. Rind 95
 C. P. Rinsland 92, 98, 100, 101
 E. J. Robbins 97
 H. M. B. Roberty 100
 A. M. Robinson 99
 J. Rogers 101
 C. Rossetti 91, 92
 J. Rostas 96
 L. S. Rothman 90, 93, 94, 96, 100,
 101, 102, 103

 M. -W. Ruf 90
 G. Russell 95

 C. P. Safvan 100
 B. C. Saha 95
 L. M. Salah 91
 H. F. Schaefer III 91
 J. Schirmer 93, 102

E. Schwartz 91
 A. Schweig 92
 N. A. Scott 96
 L. V. Shachkin 97
 R. D. Sharma 101
 V. M. Shashkov 97
 T. Yu. Sheludyakov 90
 J. M. Sichel 101
 H. Siegbahn 91
 K. Siegbahn 91
 K. J. Siemsen 99
 I. Ishii 96
 M. Simon 90
 M. A. H. Smith 92, 100
 T. H. Song 90
 V. P. Spiridonov 95
 A. T. Stair 101
 R. F. Stebbings 90
 D. Steele 99, 101
 G. Stefani 97
 T. S. Stein 100
 T. Stoecklin 89
 C. B. Suarez 101, 102
 S. Svensson 99

 M. Takekawa 89, 90
 H. Tanaka 89, 90
 M. M. Tarasenko 90
 P. R. Taylor 98
 D. T. Terwilliger 93
 J. E. Thomas 98
 F. Thommen 98
 S. X. Tian 89
 R. H. Tipping 100
 R. A. Toth 92, 100, 101
 M. Tronc 101, 102

 J. C. Ulwick 101

 A. Valentin 93, 104
 V. Valentin 95
 F. P. J. Valero 101, 102
 G. A. Vanasse 94
 J. Vander Auwera 96
 W. Vanroose 90
 K. P. Vasilevskii 102
 O. K. Voltsekhovskaya 90
 R. E. von Holdt 98
 W. von Niessen 93, 102

 S. Watanabe 89, 90
 J. W. Waters 91
 R. B. Wattson 93, 102, 103
 J. M. Weber 90
 K. H. Welge 103
 J. S. Wells 99
 D. H. Whiffen 97
 M. A. Whitehead 101
 T. A. Wiggins 100
 G. D. Willett 97
 D. Williams 91
 W. Williamson 101
 C. L. Wyatt 101

 S. Xing 103
 C. K. Xu 89
 K. Z. Xu 89

 Y. Yamaguchi 91
 S. S. Yang 90
 K. Yoshiki Franzen 103
 A. T. Young 95
 Y. L. Yung 91

 A. Zecca 89
 F. Zhang 103
 J. -P. Ziesel 89
 E. C. Zipf 94, 103
 M. Zivkovic-Rothman 90
 J. Zuniga 90, 100

Some Comments on Electron Collision Cross Sections for CO₂

The pioneer work on electron collision cross section set for CO₂ is given by J. J. Lowke in 1973. I have compiled the same set for CO₂ including new data many times. An example was shown in the book of NATO Meeting of Maratea (M. Hayashi, 1990). This cross section set is shown in Figure 1. The other cross section sets for CO₂ were presented by many authors, H. N. Kucukarpasi (1979), M. Yousfi (1987), S. E. Biagi (1991), and so on.

Now I do not like these cross section sets. Reasonable and correct cross section set have to calculate the temperature dependence of electron drift velocities at high E/N by B. Eq. and MCS method. We have solved this problem and reported at the 51th GEC Conference, Maui, as shown in this report.

Recently Michel Allan published the excellent experimental results for CO₂ + e collision data (2002), see p. 89 of this report. His used CO₂ molecules are mixture of CO₂(g), CO₂(r) and CO₂(v). The letters in the bracket mean : g is ground state, r is rotationally excited state and v is vibrationally excited state. The concentration of CO₂(g) is negligible small, and most interesting component in his experiment is the CO₂(v₁). The CO₂(v₁) of about 8 % concentration have dipole moment. Probably, concentration of (CO₂)₂ molecules are negligible small.

All molecules have the same component M(g), M(r), M(v) and M_n (n = 2), depend on the pressure and temperature condition of the experiment. Most interesting experimental results occur in the triatomic molecules. The change from linear mode combination of three atoms to bend mode combination of three atoms, or vice versa, occur very easy in the triatomic molecules by vibrational excitation.

Almost theoretical results of electron collision cross sections for molecules are for M(g). Theoretical studies involving the M(r) and M(v) are urgently required for many molecules.

I have found the interesting paper of J. R. Locker (1983)* on photoabsorption cross sections of OCS. And I would like to say that three important pioneer works on CO₂ are carried out by G. N. Haddad (1979), M. T. Elford (1980) and S. J. Buckman (1987), all in Canberra, about 20 years ago.

I would like to present our recent three conference reports at the end of this report.

M. Hayashi and Y. Nakamura : 51th GEC, Maui 265-266 (1998)

M. Hayashi and Y. Nakamura : EMS-99, Tokyo 175-176 (1999)

S. Yoshinaga, Y. Nakamura and M. Hayashi : 25th ICP1G, Nagoya 285-286 (2001)

* J. R. Locker, J. B. Burkholder, E. J. Bair and H. A. Webster III : J. Phys. Chem. 87, 1864-1868 (1983) [E, hν, OCS, 195 - 404 K at 226 nm]

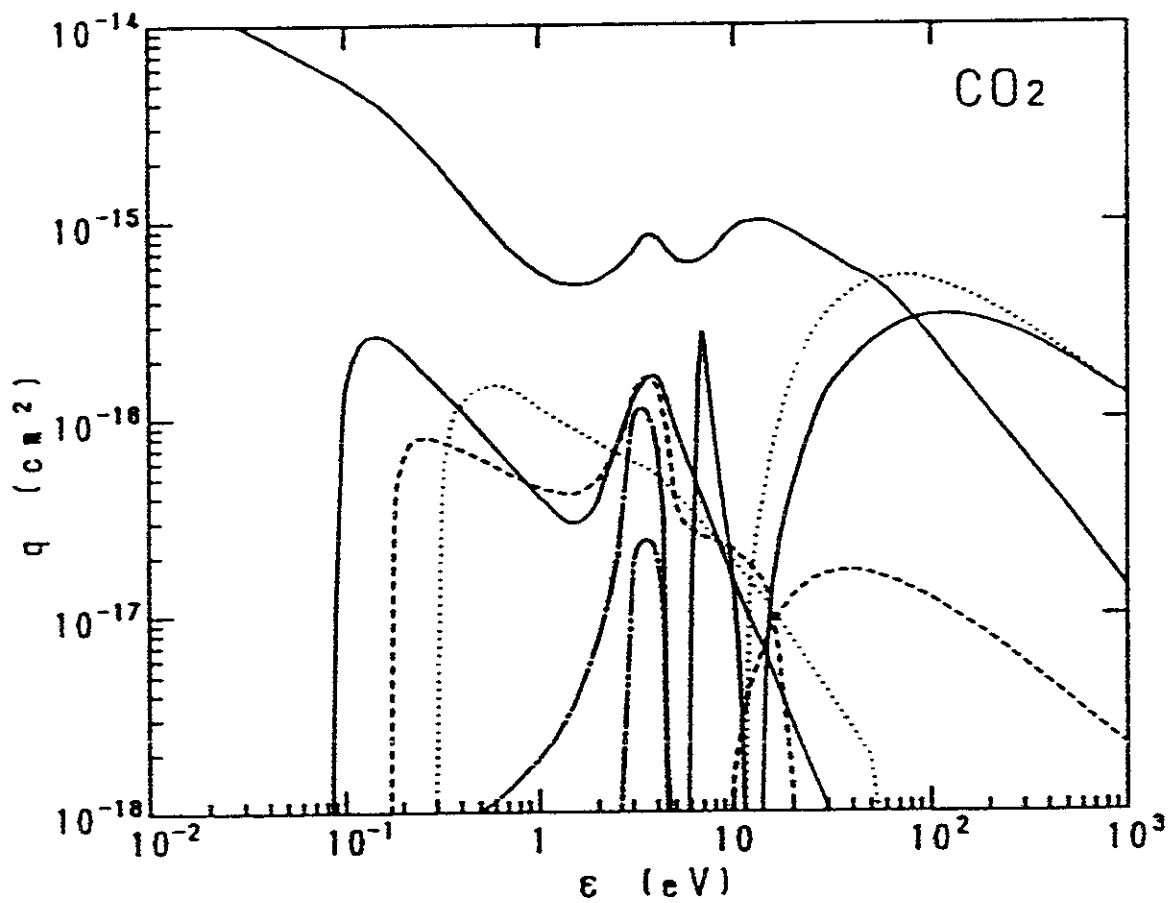


Figure 1. Electron collision cross section set for CO_2 , assuming all CO_2 molecules are in $\text{CO}_2(\text{g})$ state (M. Hayashi, 1990). In all $\text{CO}_2 + e$ experiments, CO_2 molecules are mixture of $\text{CO}_2(\text{r})$ and $\text{CO}_2(\text{v})$. So this cross section set is not applicable for exact calculations, and can use only for approximate applications.

Temperature Dependence of Electron Drift Velocity and Electron Collision Cross Section Sets for Ground State and Vibrationally Excited State of the CO₂ Molecule

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The drift velocity of electrons in carbon dioxide have been calculated at gas temperatures ranging from 193 to 573 K and at E/N values up to 100 Td assuming that the gas was a mixture of the ground state molecules and the vibrationally excited molecules and its mix ratio depended on the temperature. The calculated drift velocities agreed well with the measurement of Elford (1980).

Generally, target molecule M in beam and swarm experiments consists of the mixture of different states:

$$M = M(g) + M(r_i) + M(v_i) + M_N,$$

where M(g), M(r_i), M(v_i) and M_N represent the ground state molecules, the rotationally excited molecules, the vibrationally excited molecules and the van der Waals clusters consisting of N molecules (N≥2), respectively. It is known that the concentration of M(v_i) increases with temperature. For example, the concentration of CO₂(v₁) in CO₂ has been calculated to be 8.4 % at 313 K and 26.5 % at 673 K, respectively [1,2]. This is due to the low threshold energy (0.083 eV) of the lowest vibrational excited level (010) mode. The concentrations of rotationally excited molecules are much larger than of the vibrationally excited molecules. The concentration of the clusters are important at low temperature and high pressure conditions.

Most of theoretical calculations of electron collision cross sections have been performed for pure M(g) target, not for the M(v) molecules.

A number of different electron collision cross sections for molecules have been determined with beam and swarm experiments. So far, in most of the experiments it is implicitly assumed that the target molecules are all in the ground state, or electron collision cross section set of M(g), M(r) and M(v) are almost the same. Strictly speaking, this assumption is not correct, and in particular, cannot be applied to the molecules which have temperature dependence in electron collision cross sections.

Take the CO₂ molecule. The ground state CO₂(g) is linear, but CO₂(010) is bent and this has a dipole moment. Electron collision cross section set

of CO₂(g) and CO₂(v) are expected to be different each other. In fact, Buckman [3], Ferch [1] and Strakeljahn [4] have determined the grand total cross section Q_T for CO₂(g) and CO₂(v) and found that Q_T(v) of CO₂ are larger than Q_T(g) at the electron energies lower than about 10 eV.

It seems to be quite difficult to determine directly the electron collision cross section sets of CO₂(g) and CO₂(v) from the beam and swarm experiments. However, if the cross section sets of CO₂(g) and CO₂(v) are available, we can calculate the electron swarm parameters of a known concentration of CO₂(g) and CO₂(v) very easily. Momentum transfer cross sections q_m for CO₂(g) and CO₂(v), which have been based on the compilation of Hayashi [5] and recent experiments of Nakamura [6] and Strakeljahn [4], are shown in Fig. 1. Tentatively, we have assumed that all other inelastic cross sections of CO₂(g) and CO₂(v) are practically the same [7,8]. At a given temperature, the concentrations of CO₂(g) and CO₂(v) are known, then we have calculated the values of electron drift velocity W in CO₂. The results are shown in Fig. 2. The calculated W values at three temperatures are in general agreement with the experimental data of Elford [9].

Haddad [10] have discussed the importance of the populations of vibrationally excited CO₂ molecules already. They also have mentioned the contributions due to CO₂ molecules in rotational states and have concluded that such states are not likely to play a significant role compared with vibrationally excited molecules. But they could not reproduce their experimental data for W at high E/N from the analysis.

A similar analysis on H₂(g), H₂(r) and H₂(v)

also have to be started. The H_2 molecule have the famous long-standing controversy in the vibrational excitation cross section [11]. In the present paper, we propose to investigate and to measure the cross sections for molecules not only in the ground state but also in the excited states which should have different cross sections from those in the ground state. Although the difference among the cross section sets for $H_2(g)$, $H_2(r)$ and $H_2(v)$ may be small compared with the CO_2 molecule since the H_2 molecule is diatomic. Theoretical studies involving the excited species are urgently required.

The term "temperature dependence of cross section" may not be appropriate. The molecules, $CO_2(g)$, $CO_2(r)$ and $CO_2(v)$, have the definite and individual cross section sets, and their concentrations change with temperature. The temperature dependence is caused through different concentration of the excited components at different temperatures. Attachment cross section of the excited molecules is quite different from the ground state molecules for many attaching gases.

We wish to thank Professor H. Tawara for valuable comments.

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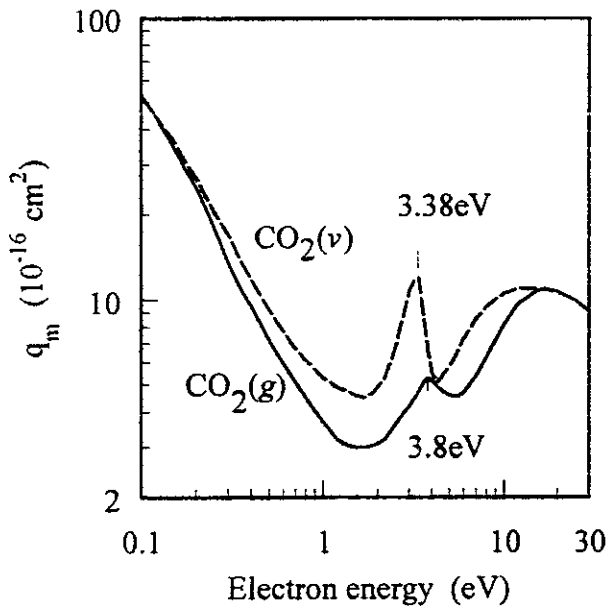


Fig. 1. The assumed elastic momentum transfer cross sections for the ground state (solid curve) and vibrationally excited CO_2 molecules (broken curve).

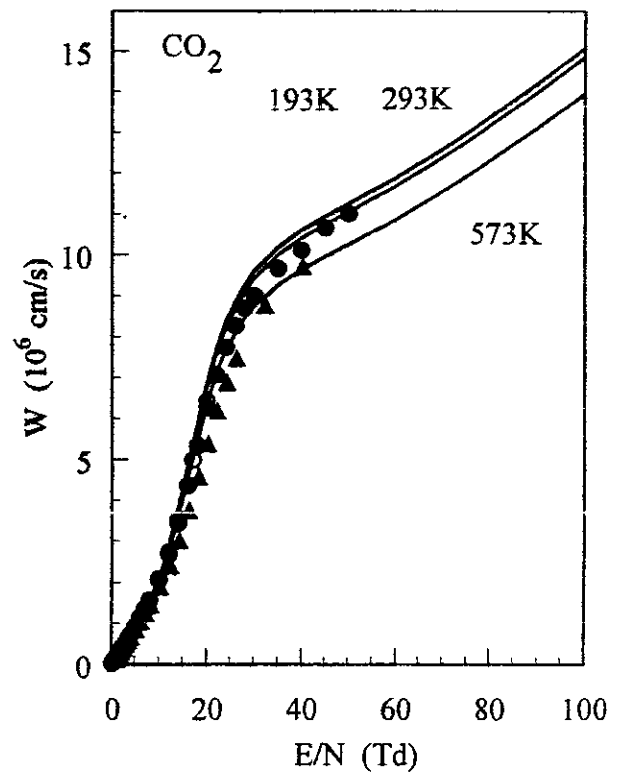


Fig. 2. Calculated electron drift velocities in CO_2 at different temperatures. Open and closed circles and closed triangles show the measurements of Elford [9] at 193, 293 and 573 K, respectively.

May We Measure the Exact Values of
Electron Collision Cross Sections
for Molecules by Beam and Swarm
Experiments ?

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We cannot measure the exact values of DCS for molecules intrinsically. Of course, we can measure the approximate values of DCS for molecules, but can measure the exact values of DCS for atoms. The reason is very simple.

Target molecule M in both beam and swarm experiments consists of the mixture of different states :

$$M = M(g) + M(r_j) + M(v_1) + M_N$$

where $M(g)$, $M(r_j)$, $M(v_1)$ and M_N represent the completely ground state molecules, the rotationally excited molecules, the vibrationally excited molecules and the van der Waals clusters consisting of N molecules ($N \geq 2$), respectively. Usually beam and swarm experiments are carried out at about 300 K. The concentration of $M(g)$ is small compared to $M(r_j)$ (except H_2 molecules) at 300 K. Then the molecule M is always mixture of $M(r_j)$ and $M(v_1)$, and the concentration of $M(r_j)$ and $M(v_1)$ changes with temperature. And electron collision cross section sets from elastic to inelastic collision processes for $M(r)$ and $M(v)$ are different each other. The target molecules M in the beam and swarm experiments are always mixtures of different molecules $M(r_j)$ and $M(v_1)$. In the case of H_2 , the target gas consists of $M(g)$ and $M(r_j)$, especially $M(g)$ and $M(r_1)$ at 78 K.

Most clear change of cross sections of $M(r_j)$ and $M(v_1)$ will occur for triatomic molecules. CO_2 , N_2O and so on. Triatomic molecules can change from linear to bend, or vice versa easily. The authors [1] have presented the different elastic momentum transfer cross sections q_{mr} and q_{mv} for $CO_2(r)$ and $CO_2(v)$, where they assumed that all other inelastic cross sections of $CO_2(r)$ and $CO_2(v)$ are practically the same. Then they have calculated the electron drift velocity W as a function of gas temperature T . When T increases, concentrations of $CO_2(v)$ increases, then W decreases with T at the same E/N , the electric field over the gas number density, around 50 Td.

We have a comment to the interesting and important paper of W. Johnstone, et al. [2]. They have measured the temperature dependence of elastic DCS for CO_2 at 4.0 eV. Unfortunately, q_{mv} for $CO_2(v)$ at 4 eV is almost equal to q_{mr} for $CO_2(r)$ [1]. Then we propose the same experiments at about 3.4 eV for CO_2 , because the temperature dependence of DCS seems to be very large there.

Winstead and McKoy [3] calculated the elastic DCS for $N_2O(g)$ at low electron energies and compare the experimental DCS data for $N_2O(r) + N_2O(v)$ mixture at 300 K. We can see large discrepancy between them at lower than about 10 eV. We urge Winstead and McKoy to calculate the DCS

for $N_2O(v)$ and also $N_2O(r)$ for comparison. We can see the same discrepancy of DCS for CO_2 [4] [5] at low electron energies.

H_2 molecules have the famous long standing controversy in the vibrational excitation cross section [6]. A possible way to solve the problem may be as follows. At first, theoreticians calculate the Q_{mg} , Q_{mr} , Q_{mv} , Q_{rg} , Q_{rr} , Q_{rv} , Q_{vg} , Q_{vr} , and Q_{vv} for $H_2(g)$, $H_2(r_1)$ and $H_2(v_1)$, as a function of electron energies. Using these data, we calculate the electron swarm parameters. From beam experiments, we cannot determine the values of Q_{mg} , Q_{mr} , Q_{rg} , Q_{rr} for $H_2(g)$ and $H_2(r_1)$ at the same time. It is clear that the threshold energies of q_r and q_v for $H_2(g)$ and $H_2(r_1)$ are different. Bhattacharyya, et al. [7] have shown that elastic integral cross sections q_{tr} are larger than q_{tg} for 20 to 200 eV for $H_2(g)$ and $H_2(r_1)$. We want the elastic DCS values for $H_2(g)$ and $H_2(r_1)$ at low electron energies lower than 10 eV. Swarm experiments also carried out in the mixtures of $H_2(g)$ and $H_2(r_1)$, except for para- $H_2(g)$ at 78 K (concentration of $H_2(g)$ is 99.3 %). We compare the experimental and calculated swarm parameters at given conditions.

Usually, theoreticians calculate the DCS for $M(g)$, not for $M(r_1)$ and $M(v_1)$, for most molecules. Theoretical studies involving the rotationally and vibrationally excited species are urgently required for many molecules. There is a interesting paper given by A. Jain [8].

If we have the cross section sets for excited inert gas clusters, we can calculate the electron swarm parameters of inert gases at high pressure and low temperature conditions. The concentration of the clusters for atoms and molecules are important at low temperature and high pressure conditions.

Most interesting temperature dependence occur for attachment cross sections q_a [9]-[13]. The values of non-dissociative q_{an} and dissociative q_{ad} are quite different for $M(r)$ and $M(v)$, and $M(r)$ and $M(v)$ have the definit and individual cross sections, independ on the temperature. Apparent temperature dependence of attachment cross section is caused through different concentration of the excited components $M(r)$ and $M(v)$ at different temperatures.

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A Measurement of Temperature dependence of Electron Transport Parameters in CO₂

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1. Introduction

CO₂ in the ground state is a linear molecule, but the vibrationally excited CO₂ has a temporal electric dipole moment and its electron collision cross section is expected to be different from that of the ground state. The CO₂ gas at elevated temperature naturally consists of not only the ground state molecules but also molecules in different vibrationally excited states and electron transport parameters are again expected to have temperature-dependence. Haddad and Elford measured the electron drift velocity in CO₂ at 193, 293 and 573K up to $E/N=40$ Td (where E is the electric field and N the gas number density, and $1\text{Td}=1\times 10^{-17}\text{ Vcm}^2$) and actually observed temperature dependent drift velocity [1]. Buckman *et al* measured the total cross section for electrons from vibrationally excited (principally 010) CO₂ molecules in the energy range 0.12-2.0 eV [2] and confirmed enhanced total cross section for the excited molecules up to about factor 2.4 of the ground state molecule.

In the present study we constructed a drift tube capable of measuring the parameters at elevated gas temperature up to about 500K, and measured the drift velocity and the longitudinal diffusion coefficient in CO₂ over the E/N range 8-300 Td at 273, 373, 423 and 473K.

2. Experimental description

Figure 1 shows the cross sectional diagram of the present electron drift apparatus. It consists of a double-shutter electron drift tube with variable drift distance, a non-inductive theathed heater-winding and a four-rod heat reflector. The drift distance is varied from 1 to 5cm by using a linear-motion feedthrough. The diameter of the drift volume is 5cm. Two pairs of thermocouples were installed at the top of the drift tube and under the electron collector and they were used to monitor and to control the inside gas temperature. The gas pressure was measured by an MKS Baratron (10 Torr Head) at cold point and the gas number density in the drift volume was determined from the monitored temperature.

Prior to actual measurement in CO₂ we measured the transport parameters in pure N₂, which are not expected to have any temperature dependence in the present E/N range, in order to confirm the determined gas number density at elevated temperatures between 293 and 473K. Figure 2 shows the measured temperature dependence of the electron drift velocity W and the product of the longitudinal diffusion coefficient and the gas number density ND_L in N₂ between 293 and 473K at several E/N .

Both parameters were almost independent on the gas temperature over the present temperature range. They also agreed very well with our previous measurement [3], and it was actually confirmed that we were able to determine the correct E/N values in the measurements.

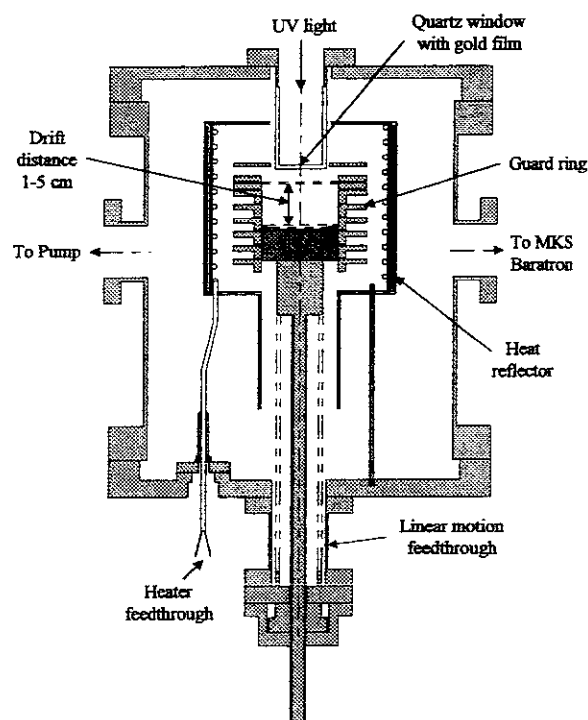


Figure 1. The cross sectional diagram of the electron drift apparatus

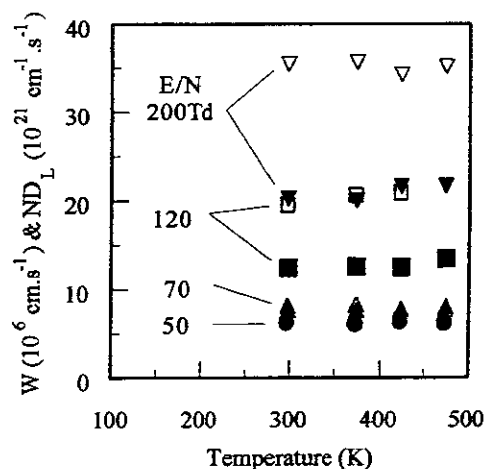


Figure 2. Temperature dependence of transport parameters in N₂. Solid and open symbols show W and ND_L , respectively.

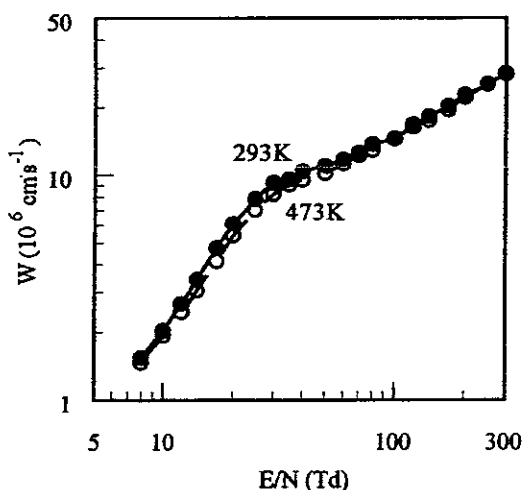


Figure 3. The electron drift velocity W versus E/N at 293K and 473K.

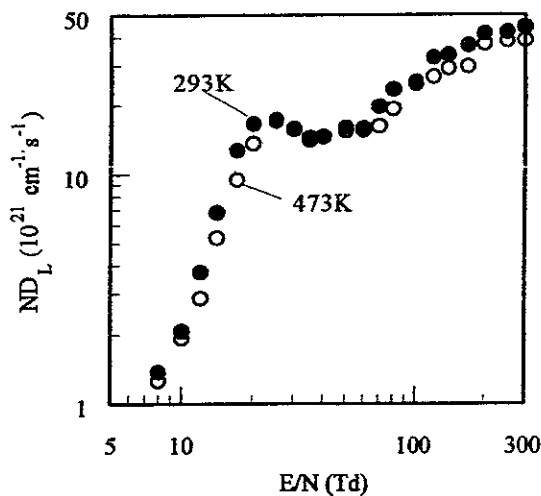


Figure 5. ND_L versus E/N at 293K and 473K.

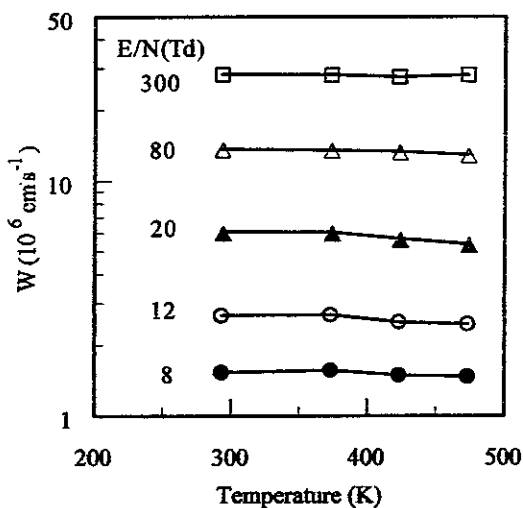


Figure 4. Temperature dependence of W in CO_2 .

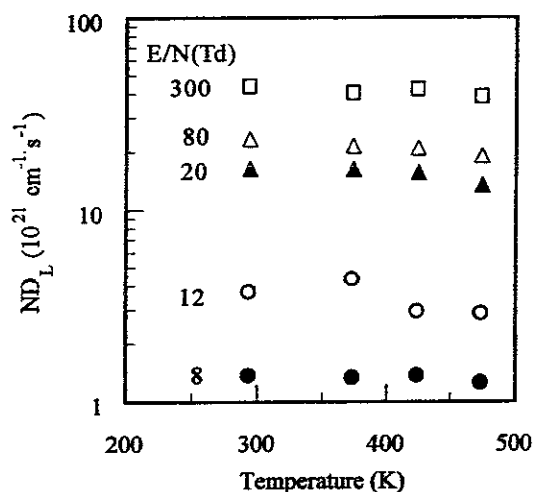


Figure 6. Temperature dependence of ND_L in CO_2 .

3. Results and discussions

Figure 3 shows the present electron drift velocities measured at 293 and 473K. At the lowest and highest E/N the drift velocity is almost independent on temperature, but it evidently decreases at higher temperature at E/N range 17-200Td.

Figure 4 shows the temperature dependence of W . The estimated error limit in W is $\pm 5\%$ and monotonic decrease in W with temperature is evident at intermediate E/N . At $E/N=8$ and 300Td W is almost independent on gas temperature.

Figure 5 shows the present ND_L at 293 and 473K, and Figure 6 shows its temperature dependence. The estimated error limit is $\pm 15\%$, but still the temperature dependence similar to the drift velocity is evident.

Mean energy of electrons participating this temperature dependence should extend up to a few eV. The dependence due to rotational excitation and deexcitation should appear even lower E/N and vibrational deexcitation should affect only very low

energy electrons. Therefore, the observed gas temperature dependence should be related to the presence of the vibrationally excited molecules with enhanced elastic cross section due to their temporal dipole moment [2].

4. Conclusions

The drift velocity and the product of the longitudinal diffusion coefficient of electrons in carbon dioxide were measured at 293, 373, 423 and 473K over the E/N range 8-300Td. A monotonic decrease with the gas temperature was observed in both transport parameters over the E/N measured except at $E/N=8$ and 300Td.

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Atoms (17)		Molecules (51)			
A + e.	A + h ν		M + e.	M + h ν .	
He 2	2170 *	2 H ₂ , D ₂	1870	5 CH ₄	750
Ne 10	1140 *	N ₂	2180		
Ar 18	1960	O ₂	1700	CF ₄	390
Kr 36	790	CO	1190	CCl ₄	210
Xe 54	940	NO	880	CCl ₂ F ₂	250
				CH ₃ Cl	90
Li 3	450	F ₂	170	SiH ₄	230
Na 11	800	Cl ₂	330	SiF ₄	140
K 19	370	Br ₂	130	GeH ₄	50
Rb 37	220	I ₂	240		
Cs 55	370			6 C ₂ H ₄	370
		HF	260	CH ₃ OH	240
O 8	390	HCl	320		
		HBr	190		
F 9	90	HI	130	7 SF ₆	780
Cl 17	130				
		3 CO ₂	1240		
				8 C ₂ H ₆	260
Cu 29	180	H ₂ O	850	C ₂ F ₆	150
Cd 48	210	O ₃	480	Si ₂ H ₆	70
Ba 56	320	N ₂ O	450		
		NO ₂	300	9 C ₃ H ₆	120
Hg 80	600	H ₂ S	270	C ₂ H ₅ OH	60
		SO ₂	260		
		CS ₂	240		
		OCS	240	11 C ₃ H ₈	190
not final, but finished mostly		4 C ₂ H ₂	390	C ₃ F ₈	100
				12 C ₄ F ₈	100
include electron swarm papers		NH ₃	400	C ₆ H ₆	240
		NF ₃	110	C ₆ F ₆	100
		BF ₃	110		
include review papers		BCl ₃	90	60 C ₆ O	300
		PH ₃	80		
		H ₂ CO	180	M _r + M _v	700

* He (Ne) + e only. Not include He (Ne) + h ν papers.

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