

§ 4. Neoclassical Transport Calculation for the N-ITB in CHS

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The parallel and radial plasma transports in multi-species plasmas in non-symmetric toroidal configurations are the important application of a recently developed neoclassical transport formulation[1]. In CHS, the toroidal flow of the fully ionized carbon in the counter direction and the reduction of the bootstrap current in the co-direction were observed when the neoclassical internal transport barrier (N-ITB) with the strong positive radial electric field ($E_r \sim 10\text{kV/m}$ at the minor radial position of $r/a \sim 0.3$) was formed in spite of the momentum input due to the co-injected heating neutral beam[2]. This phenomenon was qualitatively interpreted as the well-known mechanism making the flow in the viscosity-minimum direction that generates the counter-direction bootstrap current in the collisional regimes of the conventional helical configurations. Figure 1 shows the energy dependent pitch-angle-scattering collision frequency ν of each particle species in the particle velocity range of $v_{th}/10 < v < 5v_{th}$, where v_{th} is the thermal velocity, evaluated using the measured plasma parameters and the radial electric field at minor radial position of $r/a=0.5$ ($T_e=0.5\text{keV}$, $T_i=0.20\text{keV}$, $n_e=2.5 \times 10^{12}\text{cm}^{-3}$, $E_r=2.7\text{kV/m}$). Here, a hydrogen plasma having fully ionized carbon with the density ratio of 10% as an impurity ($Z_{eff}=3$) is assumed. The curves in Fig.1 give the integration paths to make energy-integrated viscosity and transport coefficients from the mono-energetic coefficients given as the function of ν/v and E_r/v . As shown in the figure, the collisionality of each particle species cover the wide range from the Pfirsch-Schlueter to the super banana regimes. Furthermore, the low energy components of the distributions of ions suffer the effect of so-called toroidal resonance ($E_r/B \sim vB_p/B_i$) and thus their viscosity and transport coefficients have strong dependence on the radial electric field even in the plateau regime. An connection formula is used to connect the numerically obtained $1/\nu$ regime mono-energetic diffusion coefficients to the analytical formula of the super banana regime mono-energetic diffusion coefficients[3]. The mono-energetic coefficients under the toroidal resonance can be obtained from the poloidal and toroidal viscosity coefficients in Ref.[4] using the procedure in Appendix B of Ref.[1].

Figure 2 shows the parallel flows and radial diffusion fluxes obtained for the N-ITB condition in CHS (at $r/a=0.5$) obtained by the procedure in Ref.[1]. The ambipolar condition $\sum_a e_a \Gamma_a = 0$ is attained at $E_r \sim 8\text{kV/m}$. Under this radial electric field strength, the parallel flows of the ions turn to counter direction with the velocities of 20km/s (C^{6+}) and 50km/s (H^+) and the bootstrap current $\sum_a \sum_e n_a \langle Bu_{||} \rangle / \langle B^2 \rangle^{1/2}$ in the co-direction is strongly reduced. These results qualitatively coincide with the experimental results.

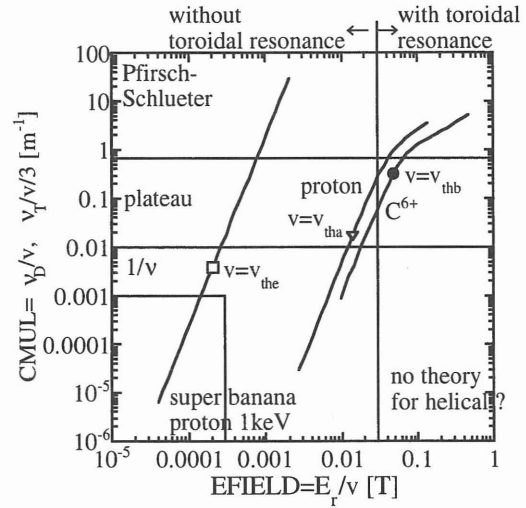


Fig.1 The energy integration path for the mono-energetic transport coefficients in ν/v - E_r/v plane. □, ▽, and ● denote the values for the thermal velocities.

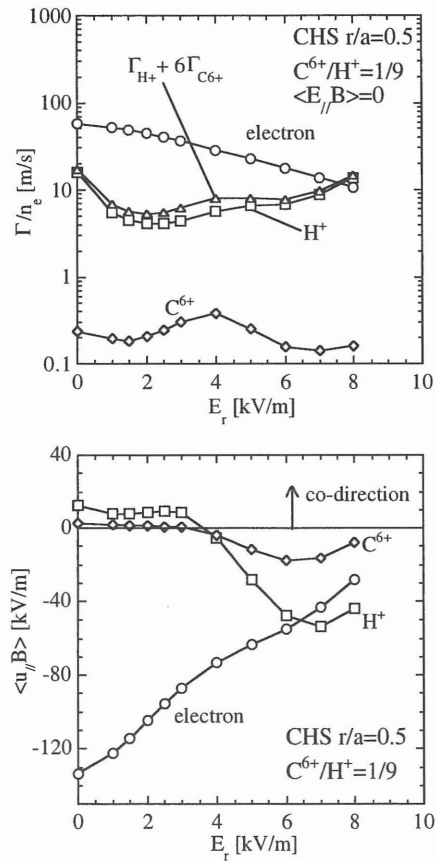


Fig.2 The parallel flows $u_{||}$ and radial diffusion fluxes Γ calculated for the N-ITB condition in CHS (at $r/a=0.5$).

References

- [1] Sugama, H. and Nishimura, S., Phys. Plasmas 9, 4637 (2002)
- [2] Ida, K., Minami, T., et al., Phys. Rev. Lett. 86, 3040 (2001)
- [3] Nishimura, S., Sugama, H., and Murakami, S., Ann. Rep. of NIFS, p.331 (Apr. 2001-Mar. 2002)
- [4] Shaing, K.C., Phys. Fluids B5, 3841 (1993)