§28. Development of a Neoclassical Moment Equation Solver II (An Application of Generalized Taguchi's Formulation as a Validation)

Nishimura, S., Sugama, H., Nakamura, Y., Nishioka, K. (Kyoto Univ.)

The traditional moment equation approach for neoclassical transport [1] is constructed by using the parallel force balance equations as $\int \mathbf{v}_{\parallel} v^{2n} d^3 \mathbf{v}$ moments of the Landau equation. This sequence of the moments naturally leads us to well-known Laguerre (Sonine) polynomial expansion of the distribution function and the gyro-phase-averaged kinetic equation. А required knowledge on the spherical harmonic order l=1 of the Coulomb collision operator is only the friction coefficient obtained by Braginskii's direct Laguerre expansion method for the operator given in the Landau's expression [2]. An important advantage of this framework is that the collisional momentum conservation and the self-adjoint property of the operator [1-4] are retained even when the polynomial expansion of the order l=1 is truncated to finite terms for converting the kinetic problem into a linear algebraic equation. However, a complicated non-diagonal coupling between the Laguerre orders may be a disadvantage. In Refs.[3,4] for e.g., the numerous non-diagonal parallel viscosity coefficients $M_{j+1, k+1}^a$ and the friction coefficients $l_{j+1, k+1}^{ab}$ with $j \neq k$ describe this coupling effect. Because of it, the obtained $\int \mathbf{v}_{\parallel} v^{2n} f_a d^3 \mathbf{v}$ moments generally depend on the choice of maximum order of the expansion j_{max} , and this number can be determined only by experiences on the convergence [3,4].

There is another weighting scheme for constructing the moment equations that may reduce the non-diagonal terms $j \neq k$ before the truncation. This second idea was originally proposed by Taguchi [5] and was generalized later in Refs.[3,6]. Even though the momentum conservation and the self-adjoint property are not guaranteed completely there, this formulation can be a useful validation method for the traditional approach. However, in spite of a fact that this second method does not require the $\int \mathbf{v}_{\parallel} v^{2n} d^3 \mathbf{v}$ sequence, all of Refs.[3,5,6] used the conventional Laguerre expansion of the Landau operator with only few terms. This is an inappropriate shortcut that should be corrected. In the present study [8], the spherical harmonic expanded RMJ (Rosenbluth-MacDonald-Judd) operator [7] without the Laguerre expansion is directly implemented in the generalized Taguchi's procedure shown in Ref.[3]. Figure 1 shows calculation examples of parallel flow moments $\langle Bu_{\parallel aj} \rangle$ of Laguerre orders $0 \le j \le 3$ given by this method in a e^{-} +H⁺ plasma at a radial position of r/a=0.5 in a magnetic

configuration with the magnetic axis position R_{ax} =3.6m in the LHD. These assumed conditions and parameters $n_e(r)$, $T_e(r)$, $T_i(r)$ are identical to those in Fig.2 in Ref.[4]. The result indicates that the generalized Taguchi's formulation [3] with the direct implementation of the RMJ operator (without the Sonine polynomial expansion) gives a good agreement with the results given by the traditional moment equation method in Ref.[4].

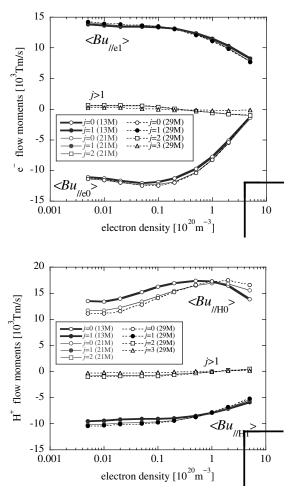


Fig.1 The $0 \le j \le 3$ Laguerre moments of the flux-surface-averaged parallel flow given by the 13M, 21M, and 29M approximations for a e^-+H^+ plasma in a LHD configuration with $R_{ax}=3.6m$, B=2.45T (at r/a=0.5).

- 1) H.Sugama and S.Nishimura, Phys.Plasmas **9**, 4637 (2002), and references cited therein.
- 2) S.I.Braginskii, Sov.Phys.JETP 6, 358 (1958).
- 3) H.Sugama and S.Nishimura, Phys.Plasmas 15, 042502 (2008).
- 4) S.Nishimura, H.Sugama, et al., Phys.Plasmas 17, 082510 (2010), 18, 069901 (2011).
- 5) M.Taguchi, Phys.Fluids B 4, 3638 (1992).
- 6) H.Maassberg, C.D.Beidler, and Y.Turkin, Phys.Plasmas 16, 072504 (2009).
- 7) I.P. Shkarofsky, et al., *The Particle Kinetics of Plasmas* (Addison-Wesley, Reading, Massachusetts, 1966) Chap.7.
- 8) S.Nishimura, et al., in ITC-20. P1-29 (Dec. 7-10, 2010)