

## §7. Effect of Gas Species on the Macroscopic Oscillations of Detach Plasma in the TPD-II

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In closed divertor configurations, neutral particles leaking through the opening of the baffle can be reduced by so called "plasma plugging effect". The effect depends on the position of the detachment front,  $z_f$ . Therefore,  $z_f$ -stability is important for the reduction of the neutrals leaking. The macroscopic  $z_f$ -oscillation around the baffle was found in the linear machine, TPD-II (Test Plasma by Direct current) [1]. In this report the effect of species of gas injected into the divertor on the  $z_f$ -oscillation is presented [2].

As shown in Fig. 1, the experimental region is partitioned into the low-pressure region (called edge plasma region: region E) and the high-pressure region (called divertor region: region D) by means of the orifice that serves as the opening of baffle plates in magnetic confinement devices. The helium plasma flows from the region E and into the region D. The upstream electron plasma density  $n_e$  is  $2 \times 10^{20} \text{ m}^{-3}$  and the electron temperature is 5 eV under the axial magnetic field of 0.25 T. The neutral gas is injected into the region D for plasma detachment, flows against the plasma flow through the orifice, and is evacuated from the region E. The experimental results were obtained under the same experimental condition except the gas species injected into the region D.

Figure 2 shows how the oscillation begins to appear, as the flow rate of injected gas,  $Q_D$ , is increased. As  $Q_D$  is increased  $P_D$  and  $P_E$  increase, and ion saturation current drawn by probe 4,  $I_{4s}$ , shown in frame (iv) decreases reflecting the movement of  $z_f$  toward the E region. At  $t \sim 350$  sec for He ( $t \sim 400$  sec for  $\text{H}_2$ ),  $z_f$  appears to reach the position of orifice. Then,  $z_f$  suddenly comes into E region, and begins to oscillate between regions of E and D. Simultaneously the values of  $P_D$ ,  $P_E$ , and  $I_{1,4}$  oscillate. It can be seen that the amplitude of the oscillation is smaller for  $\text{H}_2$  than that for He.

The detail behaviors of  $P_D$ ,  $P_E$ , and  $z_f$  during the oscillation were described in Refs. 1 and 2. The typical phenomenon, i.e., when  $z_f$  comes into the region E,  $P_D$  decreases but  $P_E$  increases, appears regardless of the gas species, which is due to the common mechanism of the plasma-plugging depending on  $z_f$ .

Figure 3 shows the  $Q_D$ -dependency of  $\Delta P$  ( $=P_D - P_E$ ). The variation of  $\Delta P$  has the upper limit,  $\Delta P_{\text{max}}$ . Whenever  $\Delta P$  reaches  $\Delta P_{\text{max}}$ ,  $z_f$  is located at the orifice. It can be considered that  $\Delta P_{\text{max}}$  balances the plasma pressure,  $P_p$ , onto the inlet of orifice:  $\Delta P_{\text{max}} \approx P_p$ . Since the amplitude of the

$\Delta P$ -oscillation should be less than the difference between  $\Delta P_{\text{max}}$  and  $\Delta P$  for without plasma indicated by dotted line in Fig. 3,  $P_p$  dominates the amplitude the oscillation. The value of  $\Delta P_{\text{max}}$  for  $\text{H}_2$  is smaller than that for He, indicating that as  $Q_D$  is increased  $P_p$  becomes smaller for  $\text{H}_2$  than that for He. This is supported by the fact depicted in Fig. 2(iv) that the decrease in the ion saturation current  $I_1$  is remarkable for  $\text{H}_2$ .

During the significant decrease in  $I_1$  in the case of  $\text{H}_2$ , Fulcher-band spectra that is a measure of the Molecular Activated Recombination (MAR) is observed from the region E. The fast recombination by MAR is able to lose  $P_p$  rapidly. Therefore, there is possibility that the small amplitude of the oscillation for  $\text{H}_2$  is owing to the MAR.

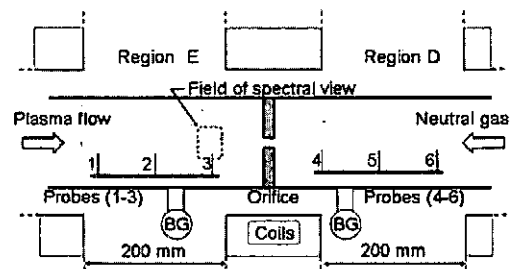


Fig.1. Schematic of experimental region of the TPD-II.

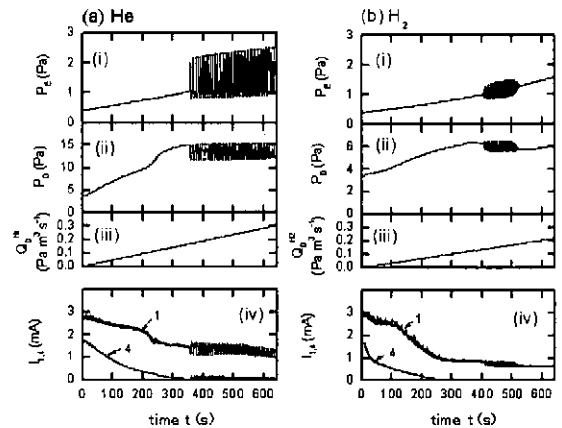


Fig.2. Long time-series showing appearance of the oscillation for cases of He (a) and  $\text{H}_2$  (b).

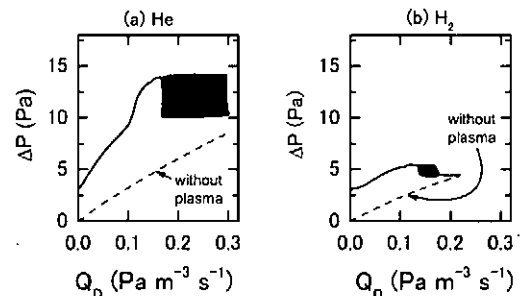


Fig.3.  $Q_D$ -dependency of  $\Delta P$  ( $=P_D - P_E$ ).

### References

- [1] Matsubara, A., et al., J. Plasma Fusion Res. 78, 196 (2002).
- [2] Matsubara, A., et al., submitted to the J. Nucl. Mater.