## §5. Favorable Aftereffect of Methane Discharges Observed in LHD Pellet Shot

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Methane (CH<sub>4</sub>) gas-puff discharge experiment has been carried out to test the controllability of an effective charge ( $Z_{eff}$ ) and the possibility of metal impurity reduction by the real time carbonization (RTC) effect. In the ideal RTC scenario, carbon atoms are effectively coated on the small wetted areas. Therefore, the total amount of CH<sub>4</sub> gas becomes much smaller than that used in the usual glow discharge method. In the experiment, the performance of the pellet shot after four discharges introducing ~20 Pa·m<sup>3</sup> of CH<sub>4</sub> was significantly improved, although the CH<sub>4</sub> discharge itself showed only slight improvement in the confinement property.

Two typical pellet shots before and after CH<sub>4</sub> discharges are compared in Fig. 1. At the time  $t = t_0 = 1.1$  s in both discharges,  $W_{\rm p}$ ,  $n_{\rm e}$ ,  $P_{\rm NB}$ , and  $T_{\rm e}$  are almost the same. However, the decay time of  $n_e$  in the pellet shot after CH<sub>4</sub> discharges (#23114) is mitigated to about twice of that before CH<sub>4</sub> discharges (#23046). This directly results in the larger plasma stored energy  $W_p$ , since the temporal behavior of the electron temperature  $T_e$  is almost the same in these discharges. Here we compare the time slice of #23046 at  $t = t_1 = 1.25$  s and that of #23114 at  $t = t_1' = 1.4$  s, to keep  $n_e$  in both slices the same. Then it can be seen that it takes 0.15 s (=  $t_1 - t_0$ ) in #23046 and 0.3 s (=  $t_1' - t_0$ ) in #23114 to obtain the same decrease in  $n_e$ . Taking into account that  $dW_p/dt \sim 0$  in these time slices, the energy confinement time in #23114 is larger than that in #23046, since  $P_{\rm NB}$  and  $n_{\rm e}$  are identical in both slices. Detailed transport analysis gives -60% reduction in the particle transport coefficient D, and -40% reduction in the electron thermal transport coefficient  $\chi_e$  at the half of the averaged minor radius. Returning to Fig. 1, the total radiation loss  $P_{\rm rad}$  and the emission of FeXVI become smaller after CH<sub>4</sub> discharges. This can be understood as the metal impurity reduction, which is expected as the real time carbonization (RTC) effect. The reduction of metal impurity irradiation can be also recognized in the soft X-ray spectrum and the radial profile of the radiation loss.

In the case discussed here, the total amount of CH<sub>4</sub> (or, the number of carbon atoms) introduced into the vacuum vessel is about 20 Pa·m<sup>3</sup> ( $\sim 5 \times 10^{21}$ ). The number of introduced carbon atoms roughly corresponds to 2 – 100 layers of graphite (the surface density of one monolayer graphite is assumed as  $\sim 3 \times 10^{19}$  atoms/m<sup>2</sup>). Here we used the area of the footprint of divertor legs ( $\sim 1 \text{ m}^2$ ) and that of the vacuum vessel directly facing the main plasma ( $\sim 300$   $m^2$ ) for the minimum- and the maximum-estimations of the wetted area. In conclusion, only a small amount of CH<sub>4</sub> gas is enough to coat the metal plasma-facing components by RTC, as was experimentally observed. The following problem is to find out the scenario, which connects the phenomenological observations of RTC effect and the improvement in both of the particle and the electron thermal transport, which cannot be explained directly by the reduction of metal impurity or the radiation loss. To answer this, more detailed and systematic experiment together with the various measurements on the electro-static fluctuation and magnetic fluctuation is needed.



Fig. 1. Waveforms of (a)  $W_p$  and  $P_{NB}$ , (b)  $n_e$  and  $P_{rad}$ , (c) the emission intensities of FeXVI and CIII, (d) the electron temperature at the plasma center  $T_{e0}$  and the pedestal ( $\rho = 0.9$ )  $T_{e_ped}$ . Broken and solid lines denote #23046 (before CH<sub>4</sub> discharges) and #23114 (after CH<sub>4</sub> discharges).